Model Studies of Outfall Systems For Desalination Plants (Part I — Flume Study)

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As the Nation's principal conservation agency, the Department of the Interior has basic responsibilities for water, fish, wildlife, mineral, land, park, and recreational resources. Indian Territorial affairs are other major concerns of America's "Department of Natural Resources".

The Department works to assure the wisest choice in managing all our resources so each will make its full contribution to a better United States—now and in the future.

FOREWORD

This is one of a continuing series of reports designed to present accounts of progress in saline water conversion and the economics of its application. Such data are expected to contribute to the long-range development of economical processes applicable to low-cost demineralization of sea and other saline water.

Except for minor editing, the data herein are as contained in a report submitted by the contractor. The data and conclusions given in the report are essentially those of the contractor and are not necessarily endorsed by the Department of the Interior.

PREFACE

The experimental investigation reported herein was conducted under Contract 14-30-2656 between the U. S. Department of the Interior and the U. S. Army Corps of Engineers. The study was conducted in the Hydraulics Division of the U. S. Army Engineer Waterways Experiment Station during the period June 1970 to July 1971 under the direction of Mr. E. P. Fortson and Mr. H. B. Simmons, Chiefs of the Hydraulics Division during this period, and Mr. T. E. Murphy, Chief of the Structures Branch. The tests were conducted by SP5 F. M. Holly, Jr., under the supervision of Mr. J. L. Grace, Jr., Chief of the Spillways and Conduits Section. This report was prepared by Messrs. Holly and Grace.

Messrs. Walter Rinne and C. L. Gransee of the Office of Saline Water, Professor R. O. Reid of Texas A&M University, and Dr. M. A. Zeitoun of Dow Chemical Company visited the Waterways Experiment Station during the investigation phase of the study to observe and discuss testing and the application of results.

COL Ernest D. Peixotto, CE, was Director of the Waterways Experiment Station during the conduct of the investigation and the preparation and publication of this report. Mr. F. R. Brown was Technical Director.

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CONVERSION FACTORS, BRITISH TO METRIC UNITS OF MEASUREMENT

British units of measurement used in this report can be converted to metric units as follows:

Multiply	By	To Obtain
inches feet square feet cubic feet feet per second cubic feet per second feet per second Fahrenheit degrees gallons (U. S.) square feet per second	2.54 0.3048 0.092903 0.02831685 0.3048 0.02831685 0.3048 5/9 3.785412 0.0930	centimeters meters square meters cubic meters meters per second* cubic meters per second meters per second per second Celsius or Kelvin degrees** cubic decimeters square meters per second

^{*} To obtain velocity in knots, multiply velocity in feet per second (fps) by 1.689.

^{**} To obtain Celsius (C) temperature readings from Fahrenheit (F) readings, use the following formula: C = (5/9)(F - 32). To obtain Kelvin (K) readings, use: K = (5/9)(F - 32) + 273.15.

A. The Problem and Purpose of Study

In the planning and design of plants for desalination of salt water, a major consideration is the environmentally acceptable disposal of the waster brine—a warm, dense, highly salt—laden effluent whose concentrations of copper and other metallic ions are considered to be a threat to the marine ecology. Among the several alternatives for disposal of this brine is the economically attractive one of discharging the effluent back into the ocean or estuary from which it was withdrawn. However, a means of mixing the dense liquid with the ambient fluid sufficiently to dilute the concentration of various salts to safe levels is required.

The Office of Saline Water has been funding an ongoing research program through the Dow Chemical Company, in which Dr. M. A. Zeitoun of Dow Chemical Company and Professor R. O. Reid of Texas A&M University have been developing conceptual designs of desalination plant outfall systems and numerical models for prediction of their performance. The purpose of the present study was to utilize a physical model to evaluate the degree of mixing attainable through use of a diffuser located on the estuary floor or the ocean floor beyond the surf zone, from which the dense brine is discharged vertically through circular ports into a uniform and steady crosscurrent.

B. Approach and Specific Objectives

Experiments at the U. S. Army Engineer Waterways Experiment Station (WES) were conducted in two areas: (1) tests of multiple-port diffusers in three distorted estuary models, to be reported under separate cover as Part II of this report, and (2) tests of single- and multiple-port diffusers at an undistorted scale of 1:20 in a flume having a level bottom and conveying uniform steady flow. The flume tests reported herein involved a study of the separate effects of the following variables on the distribution of brine downstream from a diffuser:

Variable	Prototype Range
U = ambient velocity	0.1 to 1.0 knot
V = port discharge velocity	8 to 20 fps
$\triangle \rho_{m}$ = density difference between brine and ambient fluid	0.0045 to 0.026 g/cc
D _o = port diameter	3, 6, and 9 in.

Specifically, the objectives were as follows:

1. To evaluate the effects of the above variables on the maximum

height of the upper boundary of an arcing plume, the lateral spread of the plume, and the downstream density distribution.

- 2. To determine whether single-port results can be superimposed to predict multiple-port mixing.
- 3. To determine whether or not heated brine has significantly different mixing characteristics as compared with nonheated brine.
- 4. To evaluate the mixing advantages of a multiple-port diffuser over a simple outfall pipe.

C. Qualitative Description of Jet Plume

While specific characteristics of the dense plumes will be evaluated as part of this report, it appears appropriate at the outset to describe qualitatively the general characteristics of dense jets discharged vertically into uniform ambient flow. At ambient flows only slightly above zero the jet rises nearly vertically in the longitudinal plane, arcing and falling relatively intact. Upon hitting the bottom, the brine forms rings that rapidly expand concentrically upstream and downstream in close proximity to the bottom. The effect appears to be one of gravity waves; any localized buildup of dense liquid on the bottom is unstable and must result in outward spreading to reach equilibrium.

As ambient flow is increased, the gravity wave effect is less dominant; rings form on the bottom and spread rapidly, but tend to move downstream with the ambient flow in distinct waves. At moderate ambient velocities the rings do not appear to form; the plume arcs to a peak, then flows downstream and spreads slightly as it slowly settles to the bottom.

The above discussion is descriptive of totally submerged jets. For cases in which the jet is discharged with sufficient energy to reach the surface, its characteristics are significantly altered. At low ambient velocities the plume boils and spreads concentrically along the surface; highly diluted brine then gradually falls toward the bottom. With higher ambient velocities the jet boils and spreads to an initially lesser degree than it does with low ambient velocities and is swept downstream as it spreads laterally and falls toward the bottom. However, the spread and dilution of a jet that reaches the surface are generally greater than for one totally submerged.

SECTION II: SUMMARY AND CONCLUSIONS

Nearly 400 tests were run in a 1:20-scale, uniform flow flume to evaluate the effects of port diameter, brine flow rate, density differential, and ambient velocity on the geometry and mixing characteristics of a dense jet discharged vertically through a single port. Geometry data were taken through photographic and visual observations; dilution data were compiled using combined conductivity-temperature probes. The product of the ratio of ambient to port velocities and a port densimetric Froude number has been found to be the significant parameter in all aspects of the problem.

The maximum height of the upper boundary of a jet, \mathbf{Z}_{m} , can be predicted with the following equations:

$$\frac{Z_{m} - D}{D_{o}} = C \mathbb{F}_{D}$$

$$C = 3.4 \times 10^{-0.148} (U/V_{\odot}) \mathbf{F}_{D}$$

where

D = outfall diameter, ft

 D_{\sim} = port diameter, ft

 \mathbb{E}_{D} = port densimetric Froude number

U = ambient velocity, fps

V = port velocity, fps

A correlation of the minimum dilution at a downstream station with relevant dimensionless flow parameters provides for prediction of the maximum concentrations to be expected for a given set of design/operating parameters, according to the following equation:

$$\epsilon_{\rm m} = \left[31 \times 10^{0.4 (\text{U/V}_{\rm o}) \text{ } \text{F}_{\rm D}} \right] \left(\frac{\text{x}}{\text{x}} \right)^{0.68}$$

where

x = distance downstream from diffuser, ft

x = distance at which plume falls to bottom, ft

 $\epsilon_{\rm m}$ = minimum observed dilution

Correlations of lateral plume width with downstream distance led to the following equations for prediction of plume spread:

$$\frac{W}{W_O} = \left(\frac{X}{X_O}\right)^R$$

where

w = total plume width, ft

 $w_0 = \text{plume width at } x = x_0, \text{ ft}$

and

$$R = 3.02 \times 10$$

$$-0.26(U/V_o) \mathbb{F}_D$$
for $x \le x_o$

$$R = 0.61 \log_{10} \left(\frac{U}{V_o} \mathbb{F}_D \right)$$
 for $x > x_o$

The normalizing quantities x_0 and w_0 can be predicted by

$$\mathbf{x}_{o} = 9.62 \ \mathbf{Z}_{m} \ \log_{10} \left(2 \ \frac{\mathbf{U}}{\mathbf{V}_{o}} \ \mathbf{F}_{D} \right)$$

$$\mathbf{w}_{o} = 1.51 \ \mathbf{Z}_{m} \ \log_{10} \left(4.91 \ \frac{\mathbf{U}}{\mathbf{V}_{o}} \ \mathbf{F}_{D} \right)$$

Tests using a multiple-port diffuser verified that linear superposition of single-port results can be used to predict multiple-port mixing characteristics. Tests using heated brine indicated that the presence of a temperature differential of up to 10°C between the brine and ambient fluid has no significant effect on the plume mixing characteristics. A multiple-port diffuser was found to have a significant advantage over a simple outfall pipe in keeping high concentrations of dissolved metallic ions away from the ocean floor.

SECTION III: DESCRIPTION OF TEST FACILITY

A. Flume

In choosing a scale for the laboratory model, it was important to ensure that Reynolds numbers were kept high enough so that the flows could be considered fully turbulent, as in prototype situations. A 1:20 scale of model to prototype was chosen and similitude based upon the Froudian criterion dictates the following correspondence between geometric and kinematic parameters of the two systems:

,	Prototype	Model
Length	20	ı
Area	400	1
Volume	8000	1
Time	4.4721	1 .
Velocity	4.4721	1
Discharge	1788.840	1
$\mathbb{R}_{\mathbf{p}}$	1.7×10^{5}	1.9 × 10 ³
IR _C	5.6 × 10 ⁵	6.3 × 10 ³

where

$$\mathbb{R}_{p} = \frac{V_{o}D_{o}}{v}, \mathbb{R}_{c} = \frac{UH}{v}$$

 \mathbb{R}_{n} = port Reynolds number

IR = channel Reynolds number

V = port discharge velocity

D = port diameter

v = kinematic viscosity of water

U = ambient flow velocity

H = ambient flow depth (40 ft used).

The above Reynolds numbers are representative of the minimum values simulated and indicate that all flow situations investigated were of the fully developed turbulent type. At a scale of 1:20 the model reproduces a

section of level ocean floor 140 ft wide and 600 ft long, with a maximum water depth of 40 ft. Figs. 1 and 2 are photographs of the flume. flume bottom was surrounded by 6- by 6-in, gutters that trapped dense fluid before it reflected off the flume walls. A sump area at the elevation of the gutters and a cutoff wall extending across the downstream end of the flume were built to provide an area from which excess brine could be pumped back into holding tanks. It was found during preliminary testing that it was impractical to reclaim the diluted brine, and the cutoff wall was removed. Water-surface elevations were regulated by means of a downstream gate. One wall of the flume was constructed of transparent plastic (1/2 in. and 3/4 in. thick) mounted in a wooden frame to provide for visual observations of dispersion throughout the full length of the flume. The opposite masonry wall was finished with plaster. The flume bottom was a smooth-troweled concrete slab with two coats of glossy white epoxy paint; 1.0-ft grids were painted on both vertical walls as well as on the bottom of the flume.

The experimental work reported by Dow Chemical Company verified that in modeling a dense discharge it is the density difference between the effluent and the ambient fluid, rather than the overall level of density, which is important. Therefore, the WES flume was provided with a recirculating freshwater system to model the ocean current. Fresh water supplied by pumps and a constant head tank was discharged through either a 20- or a 6-in. supply line into an 8-ft-wide forebay that was separated from the main flume by flow-straightening tiles and a rock baffle. Venturi tubes on the two supply lines provided for accurate measurement of model discharges that ranged from 0.2 to 14 cfs.

B. Brine Supply System

Two 8- by 7- by 3-ft tanks (fig. 2) were used to prepare and store brine solutions; sump pumps on the tank floor kept the solutions well mixed. Two 8-gpm centrifugal pumps with stainless steel rotors pumped from either tank through either of two Rotameters or a 1- by 1/2-in. venturi. A 2- by 2- by 1-ft tank with a point gage attached was used to calibrate volumetrically the venturi and Rotameters for model discharges ranging from 0.00014 to 0.012 cfs. The various calibrations were found to be essentially independent of the small brine density variations expected. All brine piping was 1-in. copper tubing with appropriate reducers for the pumps and Rotameters.

The model diffuser consisted of a length of pipe extending across the full width of the flume at sta 0+00. A threaded connection permitted installation of a number of different diffusers with discharge ports drilled vertically at the flume center line for most cases.

C. Flume Velocity Distribution

Velocity measurements were made in the flume to establish the degree of uniform flow obtained. Fig. 3 is a plot of velocity contours, looking

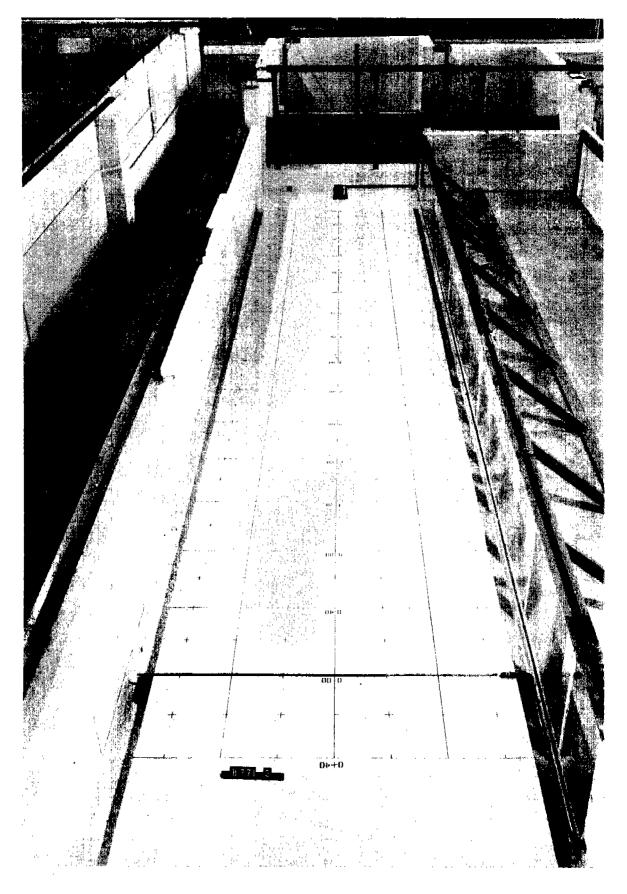


Fig. 1. Uniform flow flume (looking downstream)

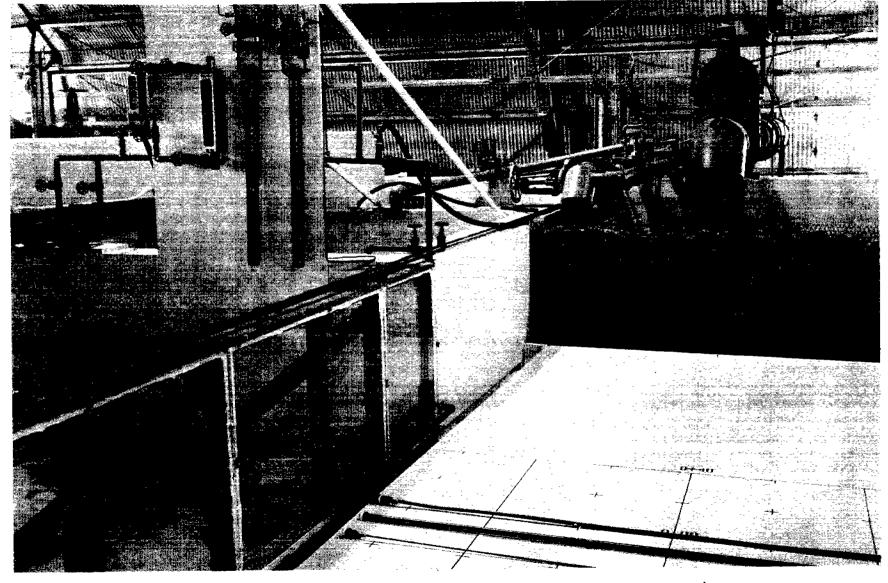


Fig. 2. Flume forebay and brine supply system (looking upstream)

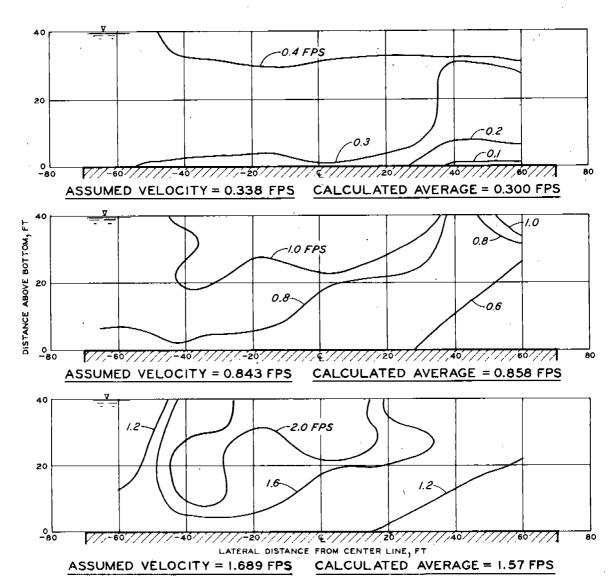


Fig. 3. Prototype velocity contours in flume (looking upstream) at x = -36 ft

upstream, determined at sta 0-36 (prototype) for assumed ambient velocities of 0.2, 0.5, and 1.0 knot (prototype). The irregular rock baffle, along with the nonsymmetrical water-supply situation, is responsible for the non-uniformity of the flow. However, the averaged velocity measurements agree well with the assumed velocities based on a discharge divided by cross-sectional area calculation; the deviations from uniform flow are not considered to be significantly different from what might be expected in a prototype ocean situation.

D. Photograph Provisions

Two 30- by 40-in. mirrors were built into a movable periscope that permitted eye-level observation of tests in the flume. Banks of photoflood lamps were placed over the flume to provide illumination for two 16-mm movie cameras that were used to photograph brine plumes through the periscope and plastic flume wall. A grid of known dimensions was placed vertically on the flume center line and photographed during initial testing for later use in scaling plume tracings.

E. Flume Instrumentation

A conductivity-temperature system was selected for use in quantifying dilution; the in situ probes were considered to have an inherent advantage over fluorescent dye methods, which require removal of a sample from the flow. A Digitec Model 501-N Digital Thermometer, made by the United Systems Corp., was used with remote probes on 50-ft leads to provide digital readout in degrees Centigrade.

The conductivity probes were designed and built at WES. Two copper electrodes were inserted into a plastic block and soldered to wire leads. The leads ran out through a length of rigid plastic tubing, which was attached to the plastic block. The entire assemblage was sealed with epoxy paint with only the electrode tips remaining bare. Each conductivity probe was then inserted into a point gage, and a thermistor probe was taped alongside it. Fig. 4 shows a typical probe assembly, nine of which were placed on rails over the flume for three-dimensional positioning within 1/2 in. of the boundaries.

A Conductivity Meter, Model R13x10-S58-P164K, made by Beckman Instruments, Inc., was used to measure the conductivity of one probe at a time. A constant resistance of 527 ohms was placed across the temperature-compensating circuit of the instrument and a 0-100 thousand ohm potentiometer was added to its bridge circuit so that the conductivity range could be varied. A 0- to 100-mv digital voltmeter and chart recorder were driven by a linearizing circuit in the conductivity meter, providing a linear record of conductivity variations. As testing proceeded it was necessary to add a so-called integrating circuit in which a capacitor accumulates voltage proportional to conductivity so that a time-averaged conductivity could be obtained.

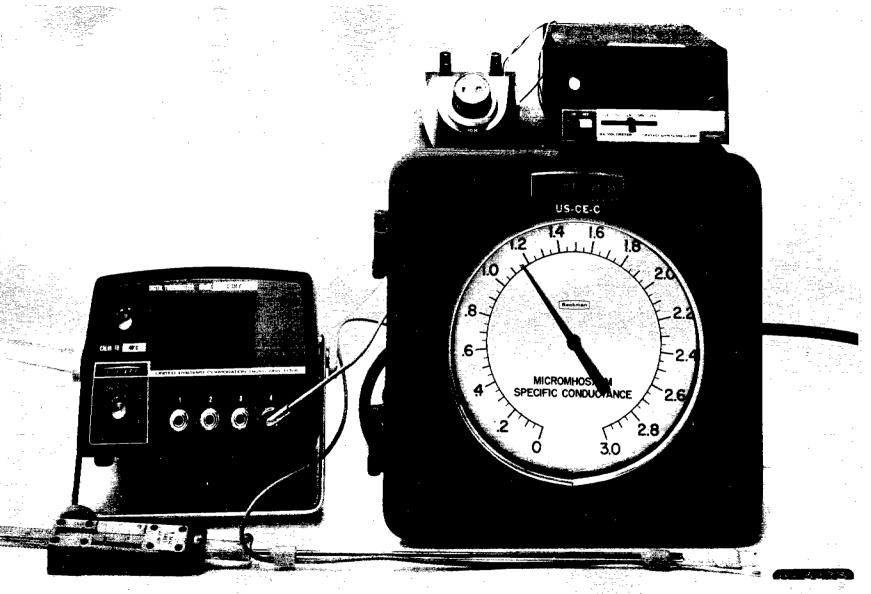


Fig. 4. Conductivity-temperature probe assembly and readout

A. General

During preliminary tests it was noted that the vertical jets tended to lean considerably in the same direction as flow in the diffuser section of the outfall. This leaning was caused by relatively high velocities in the diffuser and was reduced by increasing the outfall diameter. The final prototype outfall diameters of 20, 20, and 30 in. for the 3-, 6-, and 9-in. ports, respectively, minimized plume lean to the point that the height of the jets was not significantly affected, though the plume center line still deviated to the positive side (left side looking downstream) of the flume axis. While it was desirable to minimize plume lean in the model so that the analysis would be more straightforward, this lean could be an asset in a prototype situation where the longer plume arc length could increase overall mixing.

Brine solutions were prepared identically for both jet geometry and dilution tests. Sufficient fine-grain salt was dissolved in about 120 ft³ of fresh water to attain the desired density differential between the brine effluent and the ambient fresh water. A red food coloring was then added in sufficient quantities to give even the diluted brine a discernible color contrast with the ambient flow. After repeated density checks with hydrometers indicated that the salt had dissolved completely, the brine was pumped through the appropriate Rotameter or venturi and into the outfall and diffuser. When visual observations indicated that a steady-state condition had been reached, tests were initiated. Ambient and brine flow rates were checked frequently during tests to maintain steady-state conditions.

B. Jet Geometry Tests

Either one or two movie cameras were used to photograph the brine plume in the vertical plane through the periscope. The cameras were aimed at points 17 ft (prototype) above the flume floor, and at points 20 and 100 ft downstream. After a steady-state condition had been established in the flume, the camera(s) were turned on for approximately 10 sec. A visual sketch of the plan view spreading of the brine was then made from above the flume for about half the tests. Flow conditions were changed, and the entire procedure was repeated. Table 1 shows the test conditions for which jet geometry data were taken (see Appendix A for Notation).

C. Dilution Tests

Tests for downstream dilution were run separately from those for jet geometry, although tests having identical flow conditions were given the same number. Each series of dilution tests was preceded by a recalibration of the conductivity probes. The probes were physically prepared by filing the copper electrodes lightly to remove any surface corrosion, the buildup of which results in output signal oscillation and drift. The adjustable bridge potentiometer was set to a value that would accommodate the expected

range of conductivity. Several (three or more) calibration solutions were prepared, the first of which was pure ambient fresh water and the others were fresh water with enough brine solution added to give a range of conductivity readings up to full scale. Each probe to be used was dipped into each solution, and the temperature and conductivity were recorded. Later, the solution densities were determined on a specific gravity balance, and the corresponding temperature was again recorded. The procedure for reducing these calibrations is discussed in Appendix B.

The objective of the far-field dilution testing was to quantify the three-dimensional mixing patterns for a given operating condition from the peak of the dense plume to the downstream point where the brine spread laterally to the walls of the flume. Accordingly, conductivity-temperature probes were positioned at a number of downstream stations and detailed vertical profiles of conductivity and temperature were taken at each location with one probe at a time. Each conductivity-temperature measurement consisted of the following steps: (1) chart recorder turned on, stopwatch and integrating circuit simultaneously started; (2) temperature recorded; (3) stopwatch and integrating circuit simultaneously stopped; (4) chart recorder turned off; (5) integrating circuit voltage divided by run time and multiplied by calibration factor to get time-averaged conductivity; (6) maximum and minimum conductivity read from recorder; (7) all instruments zeroed for the next test. The ambient freshwater conductivity was also recorded for each probe as it was being used.

The basic dilution testing was conducted using a 6-in. prototype port with a density difference of about 0.021 g/cc. A few spot checks were made using an 0.01 g/cc differential, and several tests were conducted with 3-and 9-in. ports at 0.021 g/cc. A multiple-port, 20-in.-diam outfall and diffuser, with four 6-in. ports spaced at 13 ft, was tested with a density differential of 0.021 g/cc, and a 20-in.-diam simple outfall discharging horizontally with flow rates equivalent to that of the four-port diffuser was tested for comparison.

A limited test using heated brine was conducted to determine whether or not the temperature differential itself was an important factor influencing mixing. Two large space heaters were placed next to a 55-gal drum in which a brine solution was prepared. A temperature probe was installed inside a 20-in.-diam diffuser at the 6-in. port. The brine density in the heated drum was adjusted to maintain approximately an 0.021 g/cc density differential at the port. Limited downstream conductivity-temperature measurements were made, the brine temperature ranging from 4.5 to 9.4° C above the ambient temperature during the brief test.

Flow rates were checked frequently, and a visual sketch of the lateral brine spread was made for each test. Table 2 shows the dilution tests conducted. Fig. 5 presents typical plume characteristics as determined in a given test.

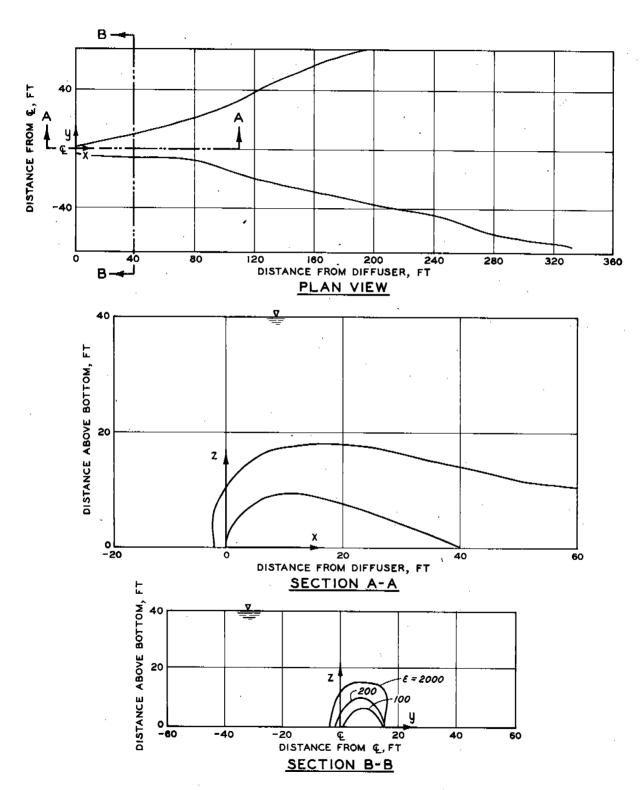


Fig. 5. Typical plume characteristics (test 639)

A. Maximum Height of Jet

The movie of a grid of known dimensions photographed at the flume center line was projected onto graph paper and the distance of the projector from the paper adjusted so that the scale of jet tracings would be 1 in. to 10 ft prototype. The scale is exact only at the center of the frame; distortion increases toward the edges of the film. Movies of colored brine plumes were projected onto graph paper, and the projector was stopped periodically so that average tracings of the shape of a plume could be made as far downstream as the color contrast permitted. A typical profile is shown in fig. 5. Appendix C (under separate cover) contains the original tracings acquired in this manner, to a scale of one major division = 10 ft prototype. Due to distortion, the apparent plume origin does not coincide with the grid origin exactly.

From each plume tracing, Z_m , the maximum height of the upper boundary of the jet, was noted and recorded (see table 1). In cases where no distinct peak was evident, the height of the upper boundary directly above the point where the lower plume boundary peaked was taken as Z_m .

During testing, it was noted that for equivalent flow conditions, the maximum height of the plume increased as the port diameter was increased. This is explained as follows. Upon its discharge from a port, the plume consists essentially of a uniform, undiluted core of constant velocity that is eroded by turbulent mixing with the ambient fluid until the turbulence has progressed all the way to the center of the plume. The undisturbed central core is being decelerated only by a modified gravity force, while the outer turbulent regions are primarily being slowed by momentum exchange with the ambient fluid. Thus, the maximum height of a plume is related to the distance required for the turbulent erosion to spread into the central core, this distance being greater for a thick jet than for a thin one.

Keffer and Baines² found in studies of a turbulent neutrally buoyant air jet perpendicular to an ambient stream that a significant parameter was the ratio of the initial jet velocity to the free stream velocity. For a dense jet it is reasonable to expect that the density difference between the jet and the ambient fluid will affect the trajectory to some degree before the plume reaches a peak, and significantly thereafter. Studies conducted by Dow Chemical Company determined that for dense jets discharging at various angles into a still fluid the normalized maximum jet height is a linear function of densimetric port Froude number. Therefore, in attempting to develop an equation for the prediction of the maximum jet height, it was assumed that

$$\frac{Z_{m}}{D_{o}} = f\left(\frac{V}{U}, T_{D}, \frac{\Delta \rho_{m}}{\rho_{f}}\right)$$
 (1)

where

Z_m = maximum height of upper boundary of the plume

D = port diameter

U = free stream velocity

 V_{O} = initial port discharge velocity $TF_{D} = \text{densimetric port Froude number}, \quad \frac{V_{O}}{\sqrt{\frac{\Delta \rho_{m}}{O_{C}}}} \text{ gD}_{O}$

 Δp_m = initial density differential between brine and ambient fluid

 ρ_f = density of ambient fluid

Correlations of (1) Zm/Do versus FD for constant Vo/U , (2) Zm/Do versus Vo/U for constant FD , and (3) Zm/Do versus Vo/U for constant $\Delta\rho_m/\rho_f$ all yielded equations for Zm/Do that satisfied most of the data but appeared to be invalid at the lower ambient velocities. This indicated that the separate effects of the dimensionless variables were being neither fully isolated nor accounted for over the entire range of flow conditions. This conclusion was confirmed by visual observations of erosion and dispersion of the jets. The dispersion in the near field appeared to be predominantly influenced by the turbulence of the jet itself or densimetric Froude number (though admittedly the pressure field and turbulence of the flowing fluid are pertinent to near-field jet dispersion), while the far-field dispersion can be primarily attributed to the relative intensity of turbulence in the far-field plume, or density current, and that of the ambient channel flow. This relative turbulence is considered to be related to the ratio of ambient to port velocities, U/V_{o} . Therefore, it was decided that analysis of data would be made in a manner such that correlation of the interrelations of U/V_{O} and F_{D} would be included in an empirical coefficient much as the frictional and form drag components are represented by an empirically determined drag coefficient.

At a conference on the study, Professor R. O. Reid of Texas A&M University indicated that Fan 3 had found the product $(\text{U/V}_{\text{O}})\mathbb{F}_{\text{D}}$ to be a significant parameter. Least-squares correlations of Z_m/D_O versus \mathbb{F}_D were made at WES for constant values of (U/Vo)FD, resulting in equations of the form

$$\frac{Z_{\rm m}}{D_{\rm o}} = C \mathbb{I}_{\rm D} + B \tag{2}$$

where B is a random intercept whose mean value is essentially zero and

$$C = \emptyset \left(\frac{V_{O}}{U} \mathbb{F}_{D} \right) \tag{3}$$

A plot of the predicted versus observed Z_m/D_O values indicated that the prediction equation was valid over the entire range of variables. The scatter was improved somewhat by accounting for the outfall diameter, D, which as reproduced in the model had the effect of elevating the entire jet a small amount. Thus Z_m/D_O was replaced by $(Z_m - D)/D_O$, and least-squares correlations of $(Z_m - D)/D_O$ versus \mathbb{F}_D for constant $(U/V_O)\mathbb{F}_D$ were repeated, leading to the following equation:

$$\frac{Z_{m} - D}{D_{C}} = C \mathbb{F}_{D} \tag{4}$$

$$-0.148(U/V_{o}) \mathbb{F}_{D}$$
 (5)

$$Z_{m} \leq H$$
 (6)

Fig. 6 is a plot of actual versus predicted values of $(Z_m - D)/D_O$.

B. Lateral Spread of Jet

Superposition of single-port dilution results to predict multiple-port characteristics requires prediction of the lateral spread of the plume. As discussed in the previous section, visual sketches of the lateral spread of the dense effluent were made for approximately half of the jet geometry tests. From these sketches, w, the observed horizontal width of the plume irrespective of the plume center line, was measured and recorded along with x, the corresponding downstream distance from the diffuser. These selected sets of coordinates, the number of which is determined by the downstream length required for the brine to spread to the edges of the flume floor for each test, are given in table 3.

The lateral spread can generally be divided into two regions: (1) before the arcing plume has settled to the bottom, and (2) after this point, when spread is generally more rapid. For a single test, w correlates linearly with x for each region as defined above. The slopes and intercepts of these correlations were related to the ambient velocity, port velocity, etc.

In attempting to develop a method for prediction of the lateral spread, three general approaches were considered and are briefly described as follows:

(1) Compute the least-squares slopes and intercepts of individual

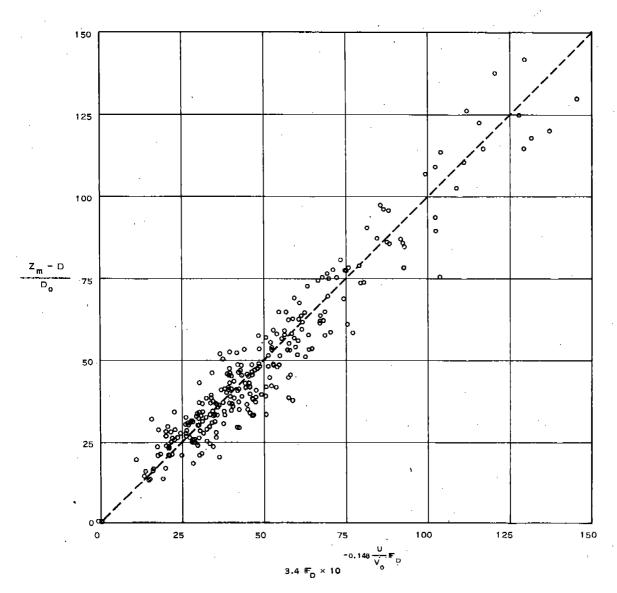


Fig. 6. Actual versus predicted values of dimensionless jet height according to equation $^{1\!\!4}$

linear correlations of w/D_O versus x/D_O for each region. Determine their functional dependence on the relevant flow parameters.

- (2) Normalize w and x by $D_o(V_O^2/U^2)$ as suggested by Keffer and Baines.² Relate $wU^2/D_oV_o^2$ to $xU^2/D_oV_o^2$ in terms of relevant flow parameters.
- (3) Defining x_O and w_O as the downstream distance at which the plume falls to the bottom and its width at that point, respectively, determine x_O and w_O as functions of \mathbb{F}_D , Z_m/D_O , and $(U/V_O)\mathbb{F}_D$. Then correlate x/x_O with w/w_O in terms of other flow parameters for both regions.

Approach 1 above yielded fair estimates of the lateral spread, but with enough systematic deviations from a perfect prediction to warrant a different approach. The success with which a modified Froude number, $(\text{U/V}_{\text{O}}) \textbf{F}_{\text{D}}$, was used to predict the maximum jet height indicated that approach 3 was worthy of consideration.

Use of this approach first requires a means of predicting \mathbf{x}_O and \mathbf{w}_O . For roughly 100 of the lateral spread sketches, \mathbf{x}_O could be approximated by the downstream distance at which the initial linear rate of spread changed to a more rapid rate. This point was often difficult to define, especially for the higher ambient velocities; all values thus determined were checked against the corresponding plume trajectory sketches (Appendix C) and a few unreasonable values revised. Values of \mathbf{w}_O , the total width of the plume at \mathbf{x}_O , were concurrently recorded.

Following the general approach of the maximum jet height correlations, $x_{\text{O}}/D_{\text{O}}$ was correlated with \mathbb{F}_{D} for constant values of $(\text{U}/\text{V}_{\text{O}})\mathbb{F}_{D}$. Individual log-log correlations were reasonably good, but their slopes and intercepts could not be correlated consistently with $(\text{U}/\text{V}_{\text{O}})\mathbb{F}_{D}$. A more successful correlation resulted from linear plots of $x_{\text{O}}/D_{\text{O}}$ and $Z_{\text{m}}/D_{\text{O}}$ for constant $(\text{U}/\text{V}_{\text{O}})\mathbb{F}_{D}$, where $Z_{\text{m}}/D_{\text{O}}$ has been assumed nearly equal to $(Z_{\text{m}}-D)/D_{\text{O}}$. Forcing these correlations to pass through the origin (assuming $x_{\text{O}}=0$ when $Z_{\text{m}}=0$), their slopes were found to be a logarithmic function of $(\text{U}/\text{V}_{\text{O}})\mathbb{F}_{D}$. An identical procedure was used for $w_{\text{O}}/D_{\text{O}}$, and the following empirical equations resulted:

$$\frac{x_{o}}{D_{o}} = 9.62 \log_{10} \left(2 \frac{U}{V_{o}} \mathbb{F}_{D}\right) \frac{Z_{m}}{D_{o}}$$
 (7)

$$\frac{\mathbf{w}_{o}}{\mathbf{D}_{o}} = 1.51 \log_{10} \left(4.91 \, \frac{\mathbf{U}}{\mathbf{V}_{o}} \, \mathbb{F}_{D} \right) \frac{\mathbf{Z}_{m}}{\mathbf{D}_{o}} \tag{8}$$

Having developed a means of predicting x_O and w_O , the relation between x/x_O and w/w_O could now be investigated. The coordinates of

total spread, as included in table 3, were divided by predicted values of $x_{\rm O}$ and $w_{\rm O}$ for each of the 187 tests for which sketches were made, and $x/x_{\rm O}$ was plotted against $w/w_{\rm O}$ for constant values of $(U/V_{\rm O})F_{\rm D}$. Least-squares linear fits of these log-log correlations (which by definition passed through $x/x_{\rm O}=1.0$, $w/w_{\rm O}=1.0$) clearly indicated an increase in rates of spread when $x/x_{\rm O}>1$, as was suggested by qualitative observations. The rates of spread were found to be an exponential function of $(U/V_{\rm O})F_{\rm D}$ for $x\leq x_{\rm O}$, and a logarithmic function for $x>x_{\rm O}$. Thus

$$\frac{\mathbf{w}}{\mathbf{w}_{o}} = \left(\frac{\mathbf{x}}{\mathbf{x}_{o}}\right)^{\mathbf{R}} \tag{9}$$

where

$$-0.26(U/V_{o}) \mathbb{F}_{D}$$

$$R = 3.02 \times 10$$
for $x \le x_{o}$ (10)

and

$$R = 0.61 \log_{10} \left(\frac{U}{V_0} \mathbb{F}_D \right) \quad \text{for } x > x_0$$
 (11)

Fig. 7 is a plot of observed values of w/w_{O} versus the values predicted using equations 7-11. Although there is a broad band of scatter due to errors in visually sketching the lateral spread and subjectively determining x_{O} and w_{O} , the overall trend indicates a valid prediction over the entire spread regime.

C. <u>Dilution for Single-Port Tests</u>

A convenient dimensionless representation of mixing is dilution, defined as

$$\epsilon = \Delta \rho_{\rm m} / \Delta \rho$$

where

$$\triangle \rho_{m} = \rho_{b} - \rho_{f}$$

 $\rho_{\rm b}$ = initial brine density, g/cc

 ρ_{f} = ambient density, g/cc

 ρ = density at some point in the far-field mixing region, g/cc

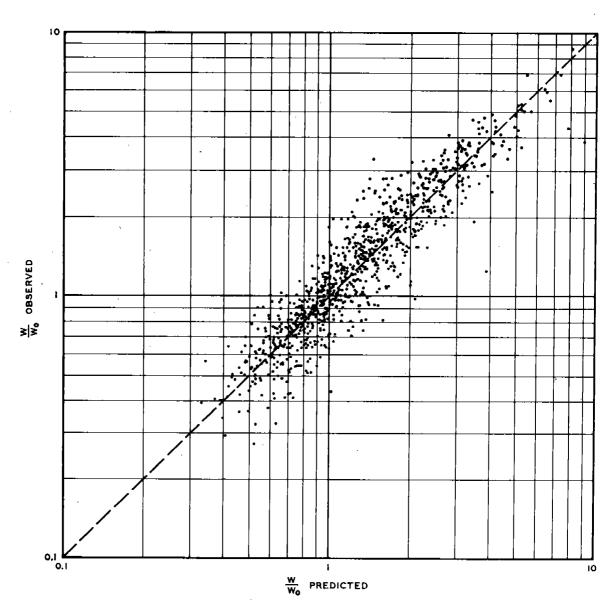


Fig. 7. Observed values of dimensionless plume width versus values predicted by equations 7-11

As defined above, the dilution of pure unmixed brine is 1.0 and that of undisturbed ambient fluid is infinity. The FORTRAN program described in Appendix B was used to calibrate the conductivity probes and to compute dilution values corresponding to minimum, average, and maximum conductivity readings at a point. Table 2 is a list of dilution test conditions. Average dilution values have been plotted at each longitudinal station for each test, and contours of constant dilution have been sketched. The resulting plots to a scale of one major division = 20 ft prototype are presented under separate cover in Appendix D. Fig. 5 is a typical sketch of dilution contours at a given cross section. The outer edges of the diluted plumes are difficult to define with precision, as the conductivity probe calibrations for salinities near the ambient salinity are extremely sensitive to slight shifts in the background conductivity. In attempting to define the limits, it should be noted that dilutions of 100, 500, and 2000 represent a 99.00, 99.80, and 99.95 percent reduction of the initial density differential, respectively.

The scope of this project prohibited any attempts to generalize the complete downstream mixing patterns. However, it was feasible to develop a prediction of the minimum dilution (i.e. maximum concentration) to be expected at any longitudinal distance from the diffuser. In general, the dilution increases with ambient velocity and downstream distance and decreases with increasing port diameter and discharge velocity.

Again, referring to the Keffer and Baines dimensionless downstream length defined as

$$\frac{xU^2}{D_0V_0^2}$$

and defining the minimum dilution as ϵ_{m} , one might expect that

$$\epsilon_{\rm m} = f\left(\frac{{\rm x}U^2}{D_{\rm o}V_{\rm o}^2}\right) \tag{12}$$

From the dilution data for each test, the minimum observed dilution ε_m , which generally was at the center of a free plume and at the bottom of brine flow along the floor, was tabulated for each longitudinal position x. These values were then plotted against $xU^2/D_0V_0^2$ on log-log axes. A general correlation was indicated, but with systematic scatter suggesting that the maximum jet height Z_m had additional bearing on the overall mixing. A second correlation using $\varepsilon_m/(Z_m/D_0)$ in place of ε_m diminished the systematic scatter, and a final log-log plot of

$$\frac{\epsilon_{\rm m}}{(Z_{\rm m}/D_{\rm o})^2} \text{ versus } \frac{xU^2}{D_{\rm o}V_{\rm o}^2}$$
 (13)

displayed good overall correlation with random scatter (fig. 8). The data for ambient velocity of 0.1 knot tended to correlate separately from the data for all other ambient velocities. The scatter in fig. 8 is due to small errors in the calibration and data reduction procedure as well as to the likelihood of "missing" a true minimum dilution which fell between discrete vertical or horizontal sampling points.

The above correlation was made before the potential of using $(\text{U/V}_{\text{O}})\mathbb{F}_{D}$ to predict jet geometry characteristics had been fully realized. Thus a new attempt to correlate the data was made along the lines of the maximum jet height and lateral spread approaches. Dilution is a measure of the degree to which ambient fluid is entrained into the brine plume. This entrainment is also the mechanism by which the plume spreads, increasing its total cross-sectional area and effective discharge. Thus the dilution should be a direct function of the rate of increase of plume area, as well as w , a parameter characteristic of the cross-sectional area. Since w/w_o correlated directly with x/x_o for constant values of $(\text{U/V}_{\text{O}})\mathbb{F}_{\text{D}}$, an obvious approach is to correlate the minimum dilution, ϵ_{m} , with x/x_o for constant values of $(\text{U/V}_{\text{O}})\mathbb{F}_{\text{D}}$.

Correlations made in this manner indicated that the log-log slopes were not a function of $(U/V_{\rm O})F_{\rm D}$, but appeared to be a single constant value for $x/x_{\rm O} \le 1.0$ and $x/x_{\rm O} > 1.0$. The log-log intercepts vary exponentially with $(U/V_{\rm O})F_{\rm D}$, and the resulting equation is as follows:

$$\epsilon_{\rm m} = \left[31.0 \times 10^{0.4 (\text{U/V}_{\odot}) \text{Fp}}\right] \left(\frac{\text{x}}{\text{x}_{\odot}}\right)^{0.68} \tag{14}$$

Fig. 9 is a plot of observed values of ϵ_m versus the corresponding predictions using equation 14. The prediction is quite good considering the difficulties in obtaining good dilution data and probe calibrations, and appears to be valid for the entire range of ambient velocity.

D. Multiple-Port Dilution Comparison

The data reduction procedures described above for the single-port dilution tests were applied to the five multiple-port diffuser tests. The resulting dilution contour plots are presented under separate cover in Appendix D.

The primary purpose of the multiple-port tests was to determine whether or not superposition of single-port results is a valid technique for prediction of multiple-port mixing. Accordingly, the single-port results of tests 655, 659, 671, and 675 were conceptually overlaid to simulate four identical plumes spaced at 13 ft o.c.; the dilution values at a single downstream station were calculated assuming linear superposition of the separate overlapping dilution contours. Fig. 10 is a plot of the calculated and observed contours; note that for tests 671 and 675 the comparison could not be made at identical downstream stations.

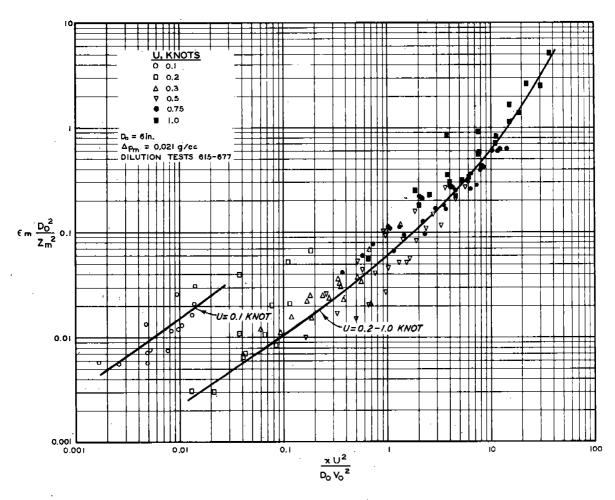


Fig. 8. Correlation of minimum dilution with dimensionless downstream distance

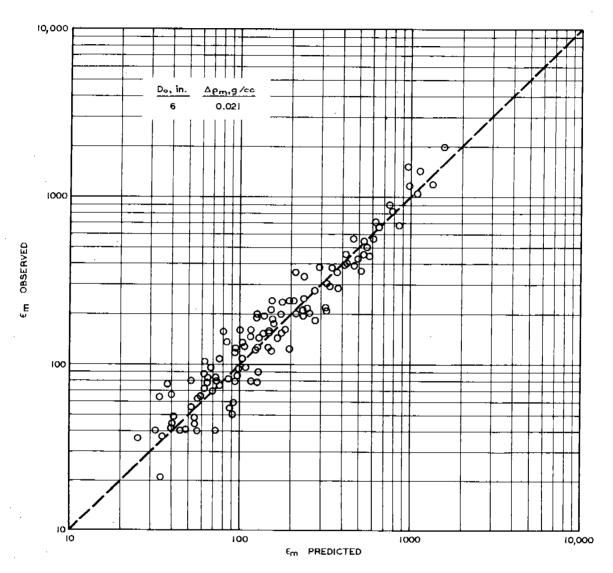


Fig. 9. Observed values of minimum dilution versus values predicted by equation 14

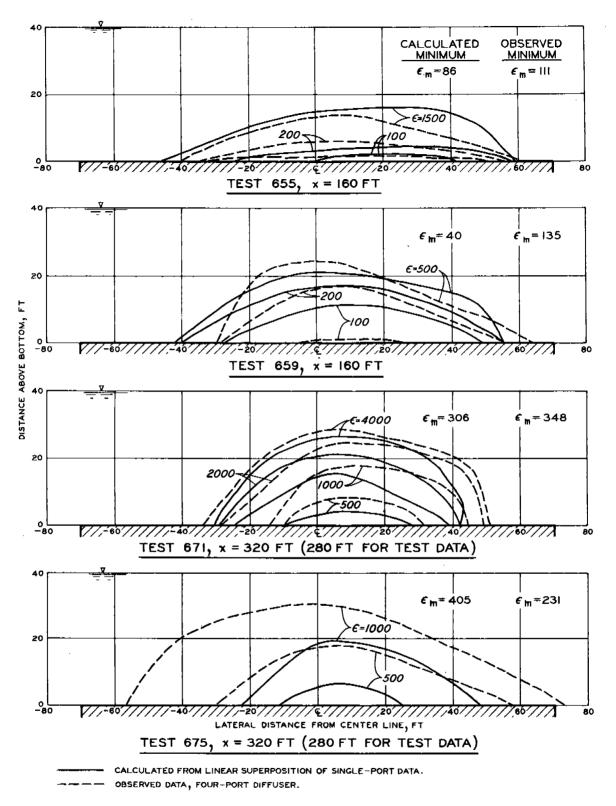


Fig. 10. Comparison of observed multiple-port dilution contours with superposition prediction (looking upstream)

The individual effects of the four separate plumes were essentially indiscernible downstream of the point where the jets merged. There appeared to be no hydrodynamic interaction between the plumes that might invalidate an assumption of linear superposition. Tests 4-655 and 4-659 agree quite well with the calculated predictions in terms of overall area of influence and degree of dilution. Tests 4-671 and 4-675 have been compared at x = 320 ft with superposition of corresponding plumes at x = 280 ft; agreement is good for 4-671 but poorer for 4-675, possibly due to the sensitivity of the results to the probe calibrations. The superposition technique does appear adequate for prediction of downstream mixing patterns.

No attempt has been made as part of this study to generalize the dilution contours. If the contours could be approximated by Gaussian distributions (as was assumed by Crew 4), the minimum dilution correlation could be used to construct a series of downstream concentration profiles that could then be superimposed to predict the mixing downstream of any multipleport diffuser, assuming a level ocean floor.

E. Simple Outfall Comparison

The brine discharged horizontally from a 20-in.-diam simple outfall tended to remain in close proximity to the bottom. Therefore, it appeared most appropriate to compare the resulting dilution patterns with the corresponding multiple-port mixing on a two-dimensional basis. Figs. 11, 12, and 13 present the plan view contours of constant dilution for corresponding simple outfall and multiple-port diffuser tests.

While the above figures may not at first suggest a dramatic difference between the two schemes, it is important to recognize that the simple outfall places the highest concentrations directly onto the ocean floor; whereas, in terms of maximum concentrations, the effective point of discharge with the diffuser is located some distance \mathbf{x}_0 downstream from the outfall where the brine has undergone an initial dilution of the order of magnitude of 100 before impinging upon the bottom. Thus, a comparison of the two schemes on the basis of bottom area affected by a given concentration would demonstrate the clear advantage of the multiple-port diffuser in protecting the ocean floor environment.

The diffuser port diameters and flow rate used in these comparison tests resulted in port Froude numbers of about 13.1, a relatively low value. Had the Froude numbers been increased by reducing the port diameters, the distance \mathbf{x}_0 , and thus the initial dilution, would have been increased. An inherent advantage of the multiple-port diffuser is that the number and size of its ports can be adjusted to provide a range of initial dilutions.

F. Effect of Heated Brine

Fig. 14 is a comparison of the dilution contours for the heated brine test with those for test 657 for which all flow conditions except the temperature differential were equivalent. Agreement is quite good at x = 40 ft and 80 ft, and acceptable at x = 120 ft and 280 ft where low concentrations amplify the probe calibration error.

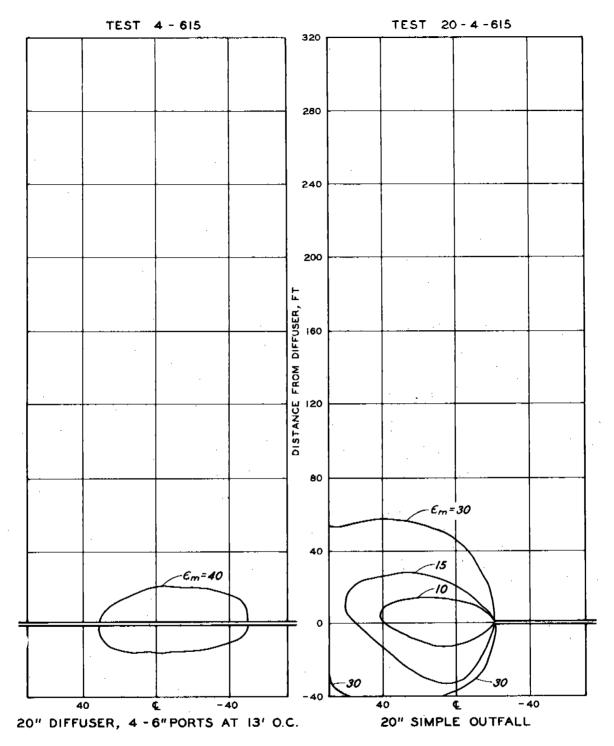


Fig. 11. Comparison of simple outfall and multiple-port diffuser; U = 0.1 knot, brine flow = 5.01 cfs, and $\Delta \rho = 0.021$ g/cc. Contours of minimum observed average dilution

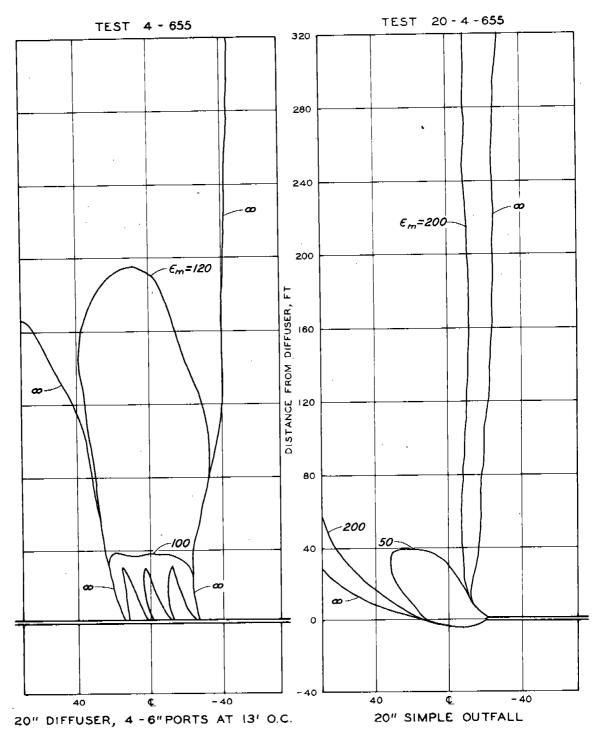


Fig. 12. Comparison of simple outfall and multiple-port diffuser; U = 0.5 knot, brine flow = 5.01 cfs, and $\Delta \rho_{\rm m}$ = 0.021 g/cc. Contours of minimum observed average dilution

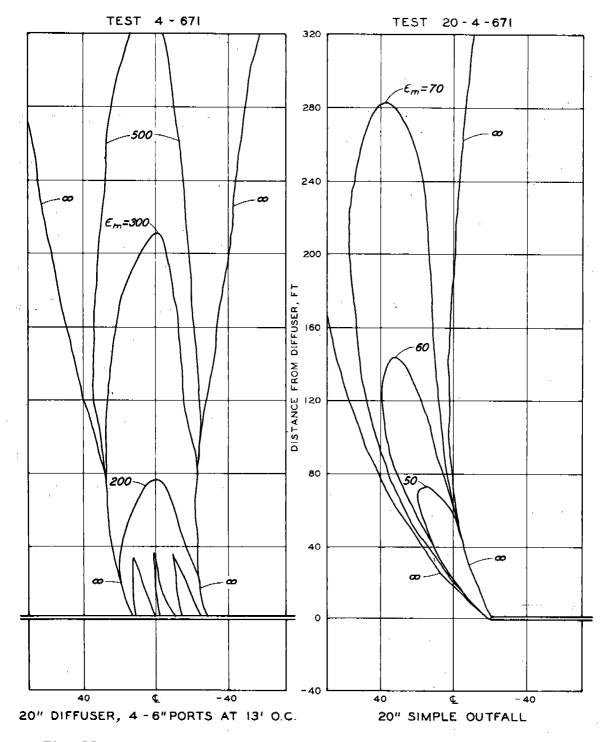


Fig. 13. Comparison of simple outfall and multiple-port diffuser; U = 1.0 knot, brine flow = 5.01 cfs, and $\Delta \rho_m$ = 0.021 g/cc. Contours of minimum observed average dilution

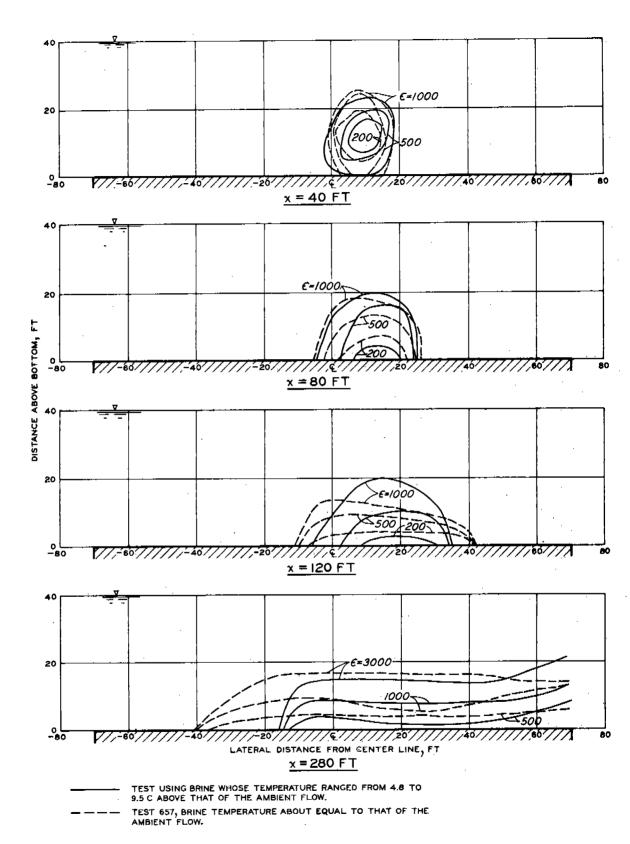


Fig. 14. Comparison of dilution contours for single 6-in. port for heated brine and test 657 (looking upstream); U = 0.5 knot, \mathbb{F}_{D} = 18.7, $\Delta \rho_{m}$ = 0.021 g/cc

A. Maximum Height of Jet

Studies at WES of submerged jets associated with lock filling and emptying systems have indicated that ratios of outfall conduit area to port area $(A_{\rm O}/A_{\rm P})$ less than 1.05 result in greatly altered manifold and jet characteristics; for higher ratios, $A_{\rm O}/A_{\rm P}$ is not considered to have an important effect on the distribution of flow in a multiple-port diffuser. In the present study, $A_{\rm O}/A_{\rm P}$ ranged from 11.1 for the 6- and 9-in. ports to 44.5 for the 3-in. ports. The existence of only two values of $A_{\rm O}/A_{\rm P}$ precluded any systematic evaluation of its effect on the jet characteristics. However, it was noted that separation of the data into two groups based on the area ratio indicated only a slight dependence on $A_{\rm O}/A_{\rm P}$, the scale of which was less than the experimental data scatter.

In a prototype situation, the level of turbulence in the ambient flow, and therefore the jet characteristics, could well be related to H , the total depth of flow. Although the few jet geometry tests run with H = 30 ft displayed no significant deviation, virtually all of the data taken here was for H=40 ft , so that all empirical constants should be considered subject to possible dependence on the depth of flow.

Although testing was not conducted below an ambient velocity of 0.1 knot prototype, setting U = 0 in equation 4 yields

$$\frac{Z_{m} - D}{D_{D}} = 3.4 \mathbb{F}_{D} \tag{15}$$

This compares quite well with the work of Turner but who predicted for a dense jet discharged vertically into a still fluid that

$$\frac{Z_{m} - D}{D_{o}} = 3.47 \, \mathbb{F}_{D} \tag{16}$$

In considering the general shape of jets, it is obvious that the mushrooming, axisymmetrical vertical jet in still fluid undergoes a transition in becoming an arcing plume at small ambient velocities. A few tests were run in hopes of determining visually at what ambient velocity this transition takes place. Although the transition point is difficult to pinpoint objectively and is to some degree a function of the densimetric Froude number, an ambient velocity of U = 0.07 knot prototype can be thought of as the transition velocity.

B. Dilution Contours

Fig. 15 is an example of the extreme fluctuations in conductivity at a point due to the turbulent jet mixing. The level of fluctuations decreased

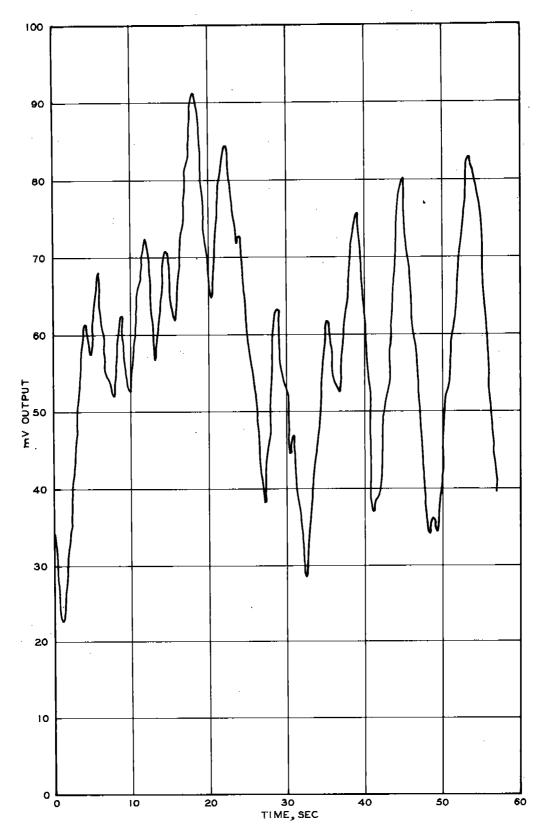


Fig. 15. Fluctuations in conductivity due to turbulent jet mixing

with distance downstream but remained significant enough that mean values could never be reliably estimated visually. As discussed earlier, 1-min samples of the conductivity probe output were processed by the integrating circuit to estimate mean values. A typical 24-min sample of conductivity output was analyzed statistically to verify the use of a 1-min sample time in actual testing.*

The overall mean of conductivities at 1-sec intervals was 61.7 mv, with a standard deviation of 10.5 mv; the total range of conductivity was from 26 to 91 mv. The sample distribution was reasonably close to a Gaussian curve with the same mean and standard deviation. Analysis of the mean values for the twenty-four 1-min samples gave a standard error of estimate of 4.5 percent. Doubling the sample time to 2 min would reduce this to 3.4 percent. The relatively small reduction in the standard error of estimate for a doubling of the sample time is considered to justify the use of a 1-min record in predicting mean values.

The dilution patterns cannot qualitatively be compared with the numerical contours predicted by Crew without evaluating the vertical turbulence exchange coefficient, a quantity that scales the dimensionless parameters used in the numerical formulation. This coefficient is a function of the scale and intensity of turbulence and was not evaluated for the WES flume. However, qualitatively the experimental contours confirm the numerical predictions of an arcing plume falling to the bottom and spreading as a gravity wave toward the flume walls as it is swept downstream.

The numerical model assumed a plane horizontal flume bottom, as was the case in the WES flume. However, tests of model diffusers in distorted estuary models at WES, reported under separate cover as Part II of this report, indicated that bottom depressions tend to fill up with relatively high concentrations of dense effluent, controlling the spread of the brine to a significant extent. Therefore, it is important to recognize that bottom irregularities in the prototype situation may cause large localized deviations from the model predictions.

C. Correlation of Dilution with Dimensionless Downstream Distance

The dilution correlation presented in fig. 9 and discussed earlier was made with data using $D_{\rm O}=6$ in. and $\Delta\rho_{m}=0.021$ g/cc. Six additional tests were run with varying values of $\Delta\rho_{m}$ and $D_{\rm O}$ and the minimum dilution values for these tests are plotted in fig. 16 for comparison with the previous dilution correlation.

Nearly all the predicted dilutions were less than the observed values. This indicates that the previously developed prediction equation, equation $1^{\rm L}$, is not valid in general, but strictly speaking can be applied only when $D_{\rm O}$ = 6 in. and $\Delta\rho_{\rm m}$ = 0.021 g/cc. However, the points plotted on fig. 16

^{*} Personal communication; analysis conducted by Professor R. O. Reid, Texas A&M University, Department of Oceanography, College Station, Tex.

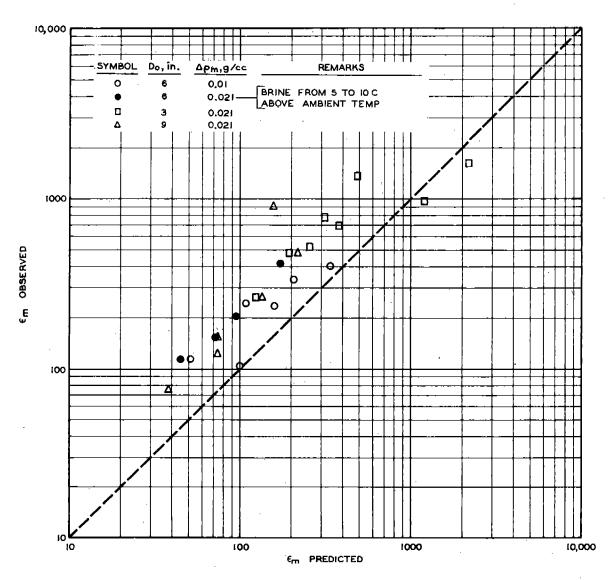


Fig. 16. Vertification of dimensionless dilution correlation, equation 14

do not fall appreciably outside the basic scatter of fig. 9, and the errors in prediction are conservative; that is, predicted dilutions may be too low. (Predicted concentrations may be too high.) This suggests that equation 14 still can be used in designing diffusers for which the concentration at some downstream point is not to exceed a specified maximum.

In developing equation 14, it was noted that, for any given value of $(\text{U/V}_{\circ})\text{F}_{\text{D}}$, the dilution is proportional to $(\text{x/x}_{\circ})^{\circ.68}$, where the exponent 0.68 is a constant for both $x/x_0 \le 1.0$ and $x/x_0 > 1.0$. On the other hand, w/w_O , the dimensionless plume width, was found to increase more rapidly for $x/x_0 > 1.0$ than for $x/x_0 \le 1.0$. This apparent contradiction results from the fact that the arcing plume for which $x/x_0 \le 1.0$ is generally round in cross section, and the width w is thus descriptive of the cross-sectional diameter, or area. On the other hand, for x/x_0 > 1.0, the plume has a more rectangular cross section, and the width alone does not fully account for the cross-sectional area. Now the entrainment of ambient fluid into the plume, which results in dilution and an increasing cross-sectional area, is governed primarily by the turbulence of the plume and the ambient fluid; therefore, the transition at $x/x_0 = 1.0$ from an arcing plume to dense flow on the bottom should not necessarily result in increased dilution or cross-sectional area, even though the rate of lateral spread does increase.

The calculated minimum dilution data for the heated brine test are also plotted on fig. 16. The systematic deviation from the prediction by equation 15 is considered to be due to the inaccurate recording of some experimental parameter, for fig. 14 demonstrates excellent agreement between this test and its nonheated counterpart. Temperature probes indicated that the heated brine reached thermal equilibrium with the ambient flow almost immediately after leaving the port. This would be less true of a jet discharged from a large port, in which a thick, undisturbed potential core would have minimum losses to the ambient fluid. From this limited test it can be tentatively concluded that a temperature differential of up to 10° C between the brine effluent and the ambient fluid will have essentially no effect on the validity of results using nonheated brine.

D. Recommended Application of Results

A primary consideration in designing desalination plant outfall systems is whether the dense plume will reach the surface, desirable for mixing but aesthetically objectionable, or remain submerged, with a decrease in overall mixing but less effect on surface appearance and recreational activity. Equation 4 provides a method of balancing port diameter, number of ports, total brine flow rate, and density differential with the ambient velocity to obtain a desired maximum jet height. Where the ambient velocity varies periodically, as in an estuary, the required rates of diversion of brine to holding tanks can be calculated for acceptable jet performance during slackwater periods.

Although generalized three-dimensional dilution patterns downstream of

a diffuser were not developed as part of this study, equation 14 does provide a means of predicting the maximum concentrations of effluent to be expected at some downstream distance from the diffuser. Thus the port diameter, number of ports, total brine flow rate, density differential, and maximum jet height can be balanced with the ambient velocity to meet established water-quality criteria downstream.

Equations 7-11 can be used to predict the lateral spread of dense plumes, which must be done before superposition of single-port results can be accomplished. As part of a proposed extended research effort, a computer program would be developed to compute three-dimensional mixing and geometry characteristics for any set of design and operation conditions, including unsteady ambient flow. Such a program would also compute specific combinations of design and operation parameters that would permit plant operation consistent with specific water-quality criteria.

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Table 1

Jet Geometry Tests

	Do		ρf	Δρ _m				7.	
Test	in.	D	g/cc	⊒m g/cc	<u>F</u> D	U Imot s	H	Z _m	ē
		in.				knots	ft	ft	
301	3	20	0.9972	0.0180	18.8	0,10	40	14,4	
302	3	20	0.9972	0.0180	24.5	0,10	40	21.0	
303	3	20	0.9972	0.0180	29.8	0,10	40	25.6	
304	3	20	0,9972	0.0180	35.7	0.10	40	30,0	•
305	3	20	0.9972	0.0180	41,4	0.10	40	36,0	
306	3	20	0.9972	0.0180	46.8	0,10	40	40.0	
307 308	3	20	0.9972	0.0180	18.8	0,21	40.	11.0	
309	3	20	0.9972	0.0180	24.5 29.8	0.21	40 40	21.8	
310	3	20 20	0.9972	0.0180	35.7	0.21	40	25.5	
311	3	20	0.9972	0.0180	41.4	0.21	40	28.8	
312	3	20	0.9972	0.0180	46.8	0,21	40	32,2	
313	3	20	0.9972	0.0180	52.4	0.21	40	37.0	
314	-3	20	0.9972	0.0182	18.4	0.31	40	12,9	
315	3	20	0.9972	0.0182	24.1	0,31	4 ŏ	13.6	
316	3	20	0.9972	0.0182	29.9	0.31	40	14.9	
317	3	20	0.9972	0.0182	35.3	0.31	40	21.0	
318	3	20	0.9972	0.0182	41.0	0.31	40	23,1	
319	3	20	0.9972	0.0182	46.4	0.31	40	28.3	
320	3	20	0.9972	0.0182	51.9	0.31	40	29.1	
321	3	20	0.997c	0.0184	18.5	0.44	40	7.9	
322	3	20	0.9970	0.0184	24:1	0,44	40	10,3	
323	3	20	0.9970	0.0184	29.7	υ,44	40	15,1	
324	3	20	0.9970	0.0184	35.1	0.44	40	17.0	
3,25	3	20	0.9970	0.0184	41.0	0.44	40	20.4	
326	3	20	0.9970	0.0184	46.2	0.44	40	24+2	-
327	3	20	0.9970	U.0184	52.0	0,44	40	23.3	
328	3	20	0.9969	0.0185	18.5	0.48	40	8,7	
329	3	20	0.9969	0.0185	24.1	0.48	40	13.1	
330	3	20	ŋ.9 9 69	0.0185	29.7	0.48	40	13.6	
331	3	20	0.9969	0.0185	35.1	0.48	40	17:2	
332 333	- 3 - 3	20	0.9969 0.9969	0.0185 0.0185	41.0 46.2	0.48	40	21.2	
334	$-\frac{3}{3}$	20	0.9969	0.0185	52.0	0,48	40	25,9	
335	3	2 ປ 2 ປ	0.9969	0.0182	18.4	0,75	40	7,5	
336	3	20	0.9969	0.0182	24.1	0.75	40	9,2	
337	3	20	0.9969	0.0182	29.9	0.75	40	8,8	
338	3	20	0,9969	0.0182	35.3	0,74	40	12,0	
339	. 3	20	0.9969	0.0182	41.0	0.75	40	13.0	
340	3	20	0.9969	0.0182	46.4	0.75	40	15.9	
341	3	20	0.9969	0.0182	51.9	0,75	40	17.8	
342	3	2 U	0.9969	U.0177	15.9	1,00	40	5,5	
343	3	20	0.9969	0.0177	24.6 30.2	1.00	40	8.8	
344	- 3	20	0.9969	0.0177		1.00	4 0	6+3	
345	3	20	0.9969	0.0177	35.7	1.00	40	9.2	
346	3	20	0.9969	0.0177	41.8	1,00	40	9.8	
347	3	20	0.9969	0.0177	47.0	1,00	40	7.5	
348	3	20	0.9969	0.0177	52.0	1,00	40	10.0	
350	3	20	0.3978	0.0105	25.U	0.10	40	20.7	
351	3	20	0.9978	0.0105 0.0105	40.3	0.10	40 40	33,1 40,0	
352 353	3	20 20	0.9978 0.9975	0.0105	55.3 70.0	0.10	40	40.0	
354	3	20 20	0.9978	0.0105	25.0	0.10	40	9.8	
355-	3 -	2 0	0.7978	0.0105	40.3	0.31	40	18.8	
	-	• ∪	0 0	0.0102		0 4 O T		T ~ 1 0	

Table 1 (Continued)

						· •••			•
	Do	D	ρ _f	$\Delta \rho_{ m m}$	_	U	н	Z _m	
Test	in.	in.	g/cc_	g/cc_	<u> </u>	knots	ft	ft	
							_		
356	3	20	0.9978	0.0105	55,3	0,31	40	25.0	
357	3	20	0.9978	0.0105	70.0	0.31	40	30.2	
358	3	2 0 ·	0.9978	0.0105	25,0	0.50	40	8,5	P. 1
359	3	20	0.9978	0.0105	40.3	0,50	40	11,4	
360	3	20	0.9978	0.0105	55,3	0,50	- 40	19.0	
361	3	20	0.9978	0.0105	70.0	0,50	40	23.0	
362	3	20	0.9978	0.0105	25.0	0.75	40	5,0	
363	3	20	0.9978	0,0105	40.3	0,75	40	7+0	
364 365	3	20	0.9978	0.0105	55.3	0,75	40	910	
365 366	3	20	0.9978	0.0105	70.0	0,75	40	12.0	
367		20	1.0002	0.0263	15.6	0.10	40-	9,8	
368	<u> 3</u> 3	<u>20</u> 20	1.0002	0.0263 0.0263	25.3 34.7	0.10	40	16.8 20.5	
369	3	20	1.0002	0.0263	44.0	0,10	40	31.0	
370	3	20	1.0002	0.0157	20.0	0,10	40	11,2	
371	3	20	1.0002	0.0157	32,3	0,10	40	22.7	
372	3	20	1.0002	0.0157	44,3	0,10	40	32.8	
373	3_	20	1.0002	0.0157	56,2	0,10	40	40.0	
374	3	20	1.0002	0.0157	20.0	0,30	40	9.0	
375	3	20	1.0002	0.0157	32,3	0,30	40	17.0	•
376	3	20	1.0002	0.0157	44.3	0.30	40	23.0	,
377	3	20	1.0002	0.0157	56,2	0,30	40	30,2	
378	3	20	1.0002	0.0263	15.6	0,30	40	6,7	,
379	3	20	1.0002	0.0263	25,3	0,30	40	11.0	
380	3	20	1.0002	0.9263	34,7	0,30	40	20.0	
381	3	20	1,0002	0.0263	44.0	0.30	40	24,0	
382	3	20	1.0002	0.0263	15.6	0,50	40	6,2	
383		20	1.0002	0.0263	25,3	0,50	40	11,5	
384	3	20	1.0002	0.0263	34,7	0.50	40	14,4	
385	3	20	1.0002	0.0263	44.0	0,50	40	20.0	
386 387	3	20 20	1.0002	0.0157	20.0 32.3	0,50 0,50	40 40	9.6	
388	3	2ù 20	1.0002	0.0157 0.0157	44.3	0,50	40	15,0 20,2	
389		20	1.0002	0.0157	36.2	0,50	40	23,4	
390	3 3	20	1.0002	0.0157	20.0	0,75	40	8,6	
391	ž	20	1.0002	0,0157	32.3	0.75	40	11,2	
392	3	20	1.0002	0.0157	44.3	0,75	40	12.0	
393	3	20	1.0002	0.0157	56,2	0,75	40	16.0	•
394	3	20	1.0002	0.0263	15.6	0,75	40	6,8	
395	3	<u> 20</u>	1.0002	0,0263	25.3	0,75	4 Q	910	
396	3	20	1.0002	0.0263	34.7	0,75	40	10:0	
397	3	20	1.0002	0,0263	44.0	0,75	40	13.0	
398	3	20	1.0002	0.0263	15.6	1.00	40	5.0	
399	3	20	1.0002	0.0263	25,3	1,00	40	8,5	
400	3	20	1.0002	0.0263	34.7 44.0	1,00 1,00	40	10:0	
401 402		20 20	1.0002	0,0263 0,0157	20.0	1,00	40	10:3 5:2	
403	3	20	1.0002	0,0157	32.3	1,00	40	8,6	
404	<u>3</u>	20	1.0002	0.0157	44,3	1.00	40	10,0	
405	3_	20	1.0002	0.0157	56.2	1,00	40	14.2	
406	3	20	1.0007	0.0045	37,0	0,10	40	21,2	
407	3.	20	1.0007	0.0045	58.0	0,10	40	34.0	
408	3	20	1.0007	0.0045	79.6	0,10	40	40.0	
409	3	20	1.0007	0.0045	37.0	0,30	40	12.0	
				(Continued)		(2	OI / S	heets)	1 100 114 115

Table 1 (Continued)

						·			•
	D	D	$^{ m P}_{ m f}$	$\Delta \rho_{ extbf{m}}$		U	Н	Zm	
Test	in.	in.	g/cc	g/cc	$\mathbf{F}_{\mathbf{D}}$	knots	ft	ft	
100	<u> </u>	<u> </u>	<u> 8</u> 7 CC	<u> </u>	<u></u>	MICOB			
410	3	20	1.0007	0.0045	58,0	0.30	40	21.3	
411	3	20	1,0007	0.0045	79.6	0,30	40	27.2	
412	<u> </u>	20	1.0007	0.0045	100.3	0,30	40	31,6	
413		20	1,0007	0.0045	37,0	0,50	40	7,8	
414		20	1.0007	0.0045	58.0	0,50	40	13,7	
415	3	20	1.0007	0.0045	79,6	0,50	40	15,8	
416	3	20	1.0007	0.0045	100.3	0,50	40	21,0	
601	6	10	0.9968	0.0192	12.5	0.10	40	21,2	
602	6	10	0.9968	0.0192	15.9	0,10	40	29,5	
603	6	10	0.9968		20.0	0,10	40	34,5	
604	6		0.9980	0.0192	24.8				
605		10		0.0192	27.0	0.10	40	40.0	
		10	0,9968	0.0192	27,3	0,10	40	40.0	
606	6	10	0.9968	0.0192	29.2	0,10	40	40.0	
607	6	10	0.9970	0.0189	12,4	0,32	40	16,6	
608	6	10	0.9970	0,0189	16,0	0,32	40	23,6	
609	6	10	0.9970	0.0189	24,9	0,32	40	26.5	
610	6	10	0.9970	0.0189	29,2	0,32	40	40.0	
611	6	10	0.9970	0.0189	13.8	0,50	40	16,2	
612	6	10	0.9970	0.0189	20.0	0,50	40	18.0	
613	6	10	0.9970	0.0189	27.4	0.50	40	25,1	
614	. 6	10	0.9970	0.0189	34.8	0,50	40	38,2	
615	6	20	0.9968	0.0190	13,7	0,10	40	27.6	
616	. 6	20	0,9968	0,0190	17.1	0,10	40	29,2	
617	- 6	20	0.9968	0.0190	19.5	0,10	40	36,1	
618	6	20	0.9968	0,0190	22.3	0,10	40	39,3	
619	6	20	0.9968	0.0190	25.8	0,10	40	40.0	
620	6	20	0.9968	0.0190	28,9	0,10	40	40,0	
621	6	20	0.9968	0.0190	33,6	0,10	40	40.0	
622	6	20	0.9968	0.0190	37.5	0,10	40	40.0	·
623	6	20	0.9968	0.0187	13.7	0,21	40	20.0	
624	6	20	0.9968	0.0187	17,1	0,21	40	25,8	
625	6	20	0.9968	0.0187	19,5	0.21	40	31,2	
626	6	2 0	0.9968	0.0187	22.3	0,21	40	32,9	•
627	6	20	0.9968	0,0187	25.8	0,21	40	40.0	
628	6	20	0.9968	0.0187	28.9	0.21	40	40.0	
629	6	20	0.9968	0.0187	33,6	0,21	40	40.0	···-
630	6	20	0.9968	0.0187	37,5	0,21	40	40.0	
631	6	20	0.9968	0.0189	13.7	0,31	40	21,2	
632	6	20	0.9968	0,0189	17.1	0.31	40	21,8	
633	6	20	0.9968	0.0189	19.5	0,31	40	25,3	
634	6	20	0.9968	0.0189	22.3	0.31	40	33,8	
635	6	20	0.9968	0.0189	25.8	0,31	40	37.9	
636	6	20	0,9968	0.0189	28.9	0,31	40	40.0	
637	6	2υ	0.9968	0.0189	33.6	0,31	40	40.0	
638	6	20	0.9968	0.0189	37.5	0.31	40	40.0	
639	6	20	0.9968	0.0185	13.8	0.44	40	18,0	
640	6	20	0.9968	0.0185	17.2	0.44	40	19,6	
641	6	20	0.9968	0.0185	19.6	0,44	40	21,8	
642	6	20	0.9968	0.0185	22.7	0.44	40	26.1	
643	- 6	20	0.9968	0.0185	26.0	0,44	40	29.8	
644	6	20	0.9968	0-0185	29,3	0.44	40	33,8	
645	6	20	0.9968	0.0185	34.0	0.44	40	40.0	
646	6	20	0.9968	0.0185	37.8	0,44	40	40.0	
647	6	20	0.9966	0.0190	13.7	0,48	40	14.3	
	=		2	(Continue			of 7 sl		
				(00000000		() (- , 61	/	

Table 1 (Continued)

		<u> </u>						
	Do	D	ρ _f	$\Delta \rho_{ m m}$	TE	U	H	Z _m
$\frac{\mathrm{Test}}{}$	in.	in.	g/cc	g/cc	<u> </u>	<u>knots</u>	$\underline{ t ft}$	<u>ft</u>
648	6	20	0.9966	0.0190	17.1	0.48	40	17.2
649	. 6	20	0.2966	0.0190	19.5	0,48	40	21,0
650	6	20	0.9966	0.0190	22.3	0,48	.40	22.5
<u>651</u>	6	20	0.9966	0.0190	25.8	0,48	40	25,9
652	6	20	0.9966	0.0190	28.9	0,48	40	32.8
653	6	<u>20</u>	0.9966	6.0190	33.6	0.48	40	34.0
654	6	20	0,9966	0.0190	37.5 13,9	0,48 0,50	40 40	40.0 14.0
655 656	6	20 20	0.9968 0.9968	0.0182 0.0182	17.3	0.50	40	17,0
657	6	20 20	0.9968	0.0182	19,9	0,50	40	20.0
658	6	20	0.9968	0.0182	22.8	0.50	40	24.1
659	6	20	0,9968	0.0182	26.2	0,50	40	28,2
660	6	20	0.9968	0.0182	29.5	0,50	4 0	30,5
661	6	2 0	<u> 0.9968</u>	0.0182	34.1	0,50	40	40.0
662	6	20	0.9968	0.0182	38.1	0,50	40	40+0
663	6	20	0,3965	0.0200	13.3	0.75	40	13.2
664	6	0	0.9965	0.0200	16.7	0,75	40	14.0
665	6	<u> 20</u>	0,9965	0.0200	19.0	0,75 0,75	40	16,6
666 667	6 6	20	0.9965 0.9965	0.0200 0.0200	21.7 25.0	0,75	40 40	20.2 23.0
668	<u>6</u>	<u>20</u> 20	0.9965	0.0200	28.1	0.75	40	22,9
669	6	20	0.9965	0.0200	32,7	0.75	40	23.9
670	- 6	20	0.9965	0.0200	36.3	0,75	40	23,9
671	6	20	0.9965	0.0198	13.3	1,00	40_	9,7
672	6	20	0.9965	0.0198	16.7	1.00	40	14.9
673	6	2 U,	0.9965	0.0198	1,9.0	1,00	40	15,9
674	6	2 0	0.9965	0.0198	21.7	1,00	40	17.8
675	6	20	0.9965	0.0198	25.0 28.1	1,00	40	23 ₁₀ 24 ₁ 6
676 677	6	2 0	0.9965 0.9965	0.0198 0.0198 ~	32.7	1.00	40	27.8
678	<u>6</u>	20 20	0.9965	0.0198	35.3	1,00	40	28,2
680	6	20	0.9983	0.0170	18.4	0.10	40	30,5
681	6	20	0.9983	5.0100	26.3	0.10	40	40.0
682	ě	2 0	0.9983	0.0100	18.4	0.21	40	24.6
683	6	20	0.9983	0.0100	26.3	0.21	40	40.0
684	6 .	20	0,9983	U • 0100	35,2	0,21	40	40.0
685	6	29	0.9983	0.0100	18.4	0,31	40	22.1
686	6	<u>2</u> ე	0.9983	0.0100	26.3	0.31	40	28,1 40,0
687	6	20 20	0.9983	0.0100 0.0100	35.2 45.8	0.31 0.31	40	40.0
688 689	6	<u>20</u>	0.9983	0.0100	18.4	0.44	40	13,8
690	- 6	20	0.9983	0.0100	26.3	0.44	40	20.9
691	6	20	0.9983	0.0100	35.2	0.44	40	30,3
692	6	20	0.9983	0.0100	45.8	0,44	4.0	40.0
693	6	2 0	0.9983	0.0100	18.4	0.50	40	15.7
694	6	20	0.9983	0.0100	26.3	0,50	40	21,0
695	6	20	0.9983	0.0100	35,2	0,50	40	28.0
696	6	20	0.9983	0.0100	45.8	0,50	40	40,0
697	6	20	0.9983	0.0100	18.4	0,75	40 40	11.9 12.0
698 699	6	20	0.9983	0.0100 0.0100	26.3 35.2	0,75	40	19,4
700	6 6	20 20	0.9963	0.0100	45.8	0.75	40	25,7
701	6	20	0.9983	0.0100	18.4	1.00	40	8.0
702	6.	20	0.9983	0.0100	26.3	1.00	40	12.0
<u> </u>	-		<u> </u>	•			<u> </u>	
				(Continued)		. (4	01 (sheets)

Table 1 (Continued)

_		<u> </u>			Λο.				7.	
		D _o	D	ρ _f	Δρ _m	<u> </u>	U .	H	Z _m	
$\frac{\mathbf{T}}{\mathbf{T}}$	est	in.	in.	g/cc	g/cc	<u> </u>	$\frac{\text{knots}}{}$	$\frac{\mathtt{ft}}{}$	<u>ft</u>	
7	03	6	20	0.998 3	0.0100	35.2	1.00	40	17.1	
	0.4	-6	20	0.9983	0.0100	45.8	1.00	40	27,5	
	05	6	20	1.0002	J.0263	11.4	0.10	40	15,5	
	06	- 6	20	1,0002	0.0263	16.4	0,10	40	18.3	
	07	6	20.	1.0002	0.0263	21.7	0.10	40	32,7	
	08	6	2 0	1.0002	0.0263	28.3	0.10	40	40.0	
	09	6	20	1.0020	0.0155	14.9	0.10	40	18.9 28.2	
	10	6	20	1.0020	0.0155 0.0155	21.5	0,10 0,10	40	40,0	
	11 12	6 · 6	20 20	1.0020	0.0155	14.9	0.20	40	19.7	
	13	6	20	1.0002	0.0155	21.5	0.20	40	29.2	
	14	6	20	1.0002	0.0155	28.5	0,20	40	40,0	
	15	6	20	1.0002	0.9263	11.4	0.20	40	40.0	
	16	- 6	20	1.0002	0.0263	16.4	0,20	40	24,3	
	17	6	20	1.0002	0.0263	21.7	0.20	4 0	33,3	
	18	6	20	1.0002	0.0263	28.3	0,20	40	40.0	
	119	6	20	1.0002	0.0263	11.4	0,50	40_	13,3	
	20	6	20	1.0002	0.0263	16.4	0.50	40	22,2	
	21	6	20	1.0002	0.0263	21.7	0.30	40	31,0	
	22	6	2 0	1,0002	0.0263	28.3 14.9	0.30	4 0 4 0	40.0 19.5	
	23	6	20	1.0002	0.0155	21.5	0,30	40	27.3	
	724 725	6	2 0	1.0002	0.0155 J.0155	28.5	0,30	40	32.6	
	26	-6	20	1.0002	0.0155	37.0	0,30	40	40.0	
	727	6	20	1.0002	0.0155	14.9	0.50	40	14,5	
	728		<u>20</u>	1.0002	0.0155	21.5	0.50	40	20,7	
	729	6	20	1.0002	0.0155	28.5	0.50	40	25,5	
	730	6	20	1.0002	0.0155	37.0	0,50	40	30.8	
	731	6	20	1.0002	0.0263	11.4	0,50	40	12,0	
	732	6	20	1.0002	0.0263	16.4	0,50	4 0	18,2	
	733	6	20	1.0002	0.0263	21.7	0,50	40	25.0	
	734	6	2 U	1.0002	0.0263	28.3	0.50	40	31,2	
	735	. <u>6</u>	20	1.0002	0.0263	11,4	0.75 0.75	40	10.0	
	736 737	6	20	1.0002 1.0002	0.0263 0.0263	21.7	0.75	40	17.0	
	738	- 3	<u>20</u> 20	1.0002	0.0263	28,3	u.75	40	21.3	
	739	6	20	1.0002	0.0153	14.9	0.75	40	12.9	
	740	6	<u>2</u> υ	1.0002	0.0153	21.5	0,75	40	16,6	
	741	6	20	1.0002	0.0153	28,5	0,75	40	20.4	
	742	6	20	1.0002	0.0153	37.0	0.75	40	22.7	
	743	6	20	1.0002	0.0153	14.9	1,00	40	17,4	
	744	6	20	1.0002	0.0153	21.5	1.00	40	18,5	
	745	<u> </u>	20	1.0002	0.0153	28,5	1,00	40	20.0	
	746 747	6	2 0 2 0	1.0002 1.0002	0.01 53 0.02 63	37.0 11.4	1,00 1,00	40 40	21,5 9,6	
	748	- 6	20	1.0002	0.0263	16.4	1,00	40	14,5	
	749	6	20	1.0002	0.0263	21.7	1,00	40	18,5	
	750	6	20	1.0002	0.0263	28.3	1,00	40	21.6	_
7	751	6	2 0	1.0005	0.0046	25.5	0,10	40	40.0	
	752	- 6	20	1.0005	0.0046	38,3	0,10	40	40.0	
	753	6	20	1,0005	0.0046	26,5	0,30	40	20+0	
	754	6	20	1.0005	0.0046	38.3	0,30	40	30.3	
	755 722	<u>6</u>	20	1.0005	0.0046	50.8	0,30	40	40+0	
•	756	6	20	1.0005	0.0046	65.7	0,30	40	40.0	
					(Continued)		(5_	of 7 s	heets)	

Table 1 (Continued)

		Do		. 0 .	Δo _m				7.	
	To at	_0	D	ρ _f		_ IF _D	Ŭ Januari e	H	Z m	,
	<u>Test</u>	in.	in.	g/cc	g/cc	<u>. D</u>	<u>knots</u>	<u>ft</u>	<u>ft</u>	
_	757	6	20	1.0005	0.0046	26.5	0.50	40	18,6	
	758	6	20	1,0005	0.0046	38.3	0,50	40	19.8	.
	759	6	20	1.0005	0.0046	50.8	0,50	40	28.6	• ,
	760 761	6	20	1.0005	0.0046	26.5	0,50 0,75	40	30,8 12,2	
	762	6	20	1.0005	0.0046	38,3	0.75	40	15.0	*
	763	6	20	1.0005	0.0046	50.8	0,75	40	20.2	
	764	6	2 0	1.0005	0.0046	65.7	0,75	40	23.0	
	765	6	20	1.0005	0.0046	26,5	1,00	40	1113	" -
	766	6	<u> 2) </u>	1.0005	0.0046	50.8	1,00	4.0	16,4	
	901	9	30	0.9967	0.0198	11.9	0.10	40	33.1	
	902	9	30 30	0.9967	0.0198 0.0198	15.1 18.0	0,10	40	40.0	
	904	9	30	0.9967	0.0198	21.1	0.10	40	40.0	
	905	- 6 -	3 0	0.9967	0.0198	23.5	0,10	40	40.0	
	906	9	30.	0.9967	0.0198	26.4	0.10	4:0	40.0	
	907	9	30	0.9967	C.0198	29.0	0.10	40	40.0	
	908	9	30	0.9967	0,0196	11.9	0,21	40	31.0	
	909	9	30	0.9967	0.0196	15.1	0,21	40	37.3	
-	910 911	9	30 30	0.9967 0.9967	0.0196 0.0196	18.0 21.1	0,21	40	40.0	
	912	9	3 0	0.9967	0.0196	23.5	0.21	40	40.0	
	913	- 9 -	30	0.9967	0.0196	26.4	0,21	40	40.0	
	914	9	30	0.9967	0.0196	29.0	0,21	40	40.0	
	915	9	30	0.9967	0.0196	11.9	0,31	40	29,8	•
	916	9	3 0	0.9967	0.0196	15.1	0,31	40	37.8	
	917	9	30	0.9967	0.0196	18.0	0.31 0.31	40 40	40.0 40.0	1
	918 919	9	3 0	0.9967 0.9967	0.0196	21.1 23.5	0,31	40	40.0	
	920	9	3 0	0.9967	0.0196	26.4	0.31	40	40.0	·
	921	.:9	30	0.9967	0.0196	29.0	0.31	40	40.0	
	922	9	3 0	0.9963	0.0203	<u>_11,8</u>	0.44	40	26.0	·
-	923	9	3 U	0.9963	0.0203	15.0	0.44	40	29.7	•
	924	9	30	0,9963	0.0203	17,9 21,0	0,44	40	38,8 40.0	· · · · · · ·
	925 926	9	3 υ 30	0.9963	0.0203 0.0203	23,3	0.44	40	40.0	
, .	927	- 6	″ 3 ນ	0.9963	0.0203	26.2	0.44	40	40.0	
	928	9	30	0.9963	0.0203	28.7	0,44	40	40.0	
-	930	9	3 0	0.9966	0.0199	11.9	0.48	40	24.6	
	931	<u>9</u>	<u> 30</u>	0.9966	0.0199	15.1	0.48	40	2812	
	932 933	9	3 0	0.9966 0.9966	0.0199 0.0199	18.0 21.1	0.48 0.48	4 0 4 0	35.0 40.0	
	934	9	<u> 3ე</u> 3ე	0.9966	0.0199	23.5	U.48	40	40.0	
	935	9	<u> 3</u> บั	0.9966	0,0199	26.4	0,48	40_	40.0	
	936	9	30	0.9966	0.0199	29.0	0,48	40	40.0	
	937	<u> </u>	30	0,9967	0.0197	11.9	0.50	40	25.1	
	938	9	30	0.9967	0.0197	15.1	0.50	40	29.9	
	939	9	30 30	0.9967	0.0197 0.0197	18.0 21.1	0,50	4·0 4·0	36,2 40,0	
•	941	9	30_	0,9967	0.0197	23.5	0,50	40	40.0	
	942	9	30	0.9997	0.0197	26,4	0.50	40	40.0	
	943	9	3.0	0.9967	0.0197	29,0	0,50	40	40.0	
	944	9	3 0	0.9968	0.0196	11.9	0,75	40	21:1	
	945	9	3 0	0,9968	0.0196	15.1	0,75	40_	22,0	
		,	Α		(Continued)		(6	of 7 s	heets)	,

Table 1 (Concluded)

	Do	D	ρ _f	$\Delta \rho_{\underline{m}}$	_	υ .	H	Z _m
Test	in.	in.	g/cc	g/cc	F _D	knots	ft	<u>ft</u>
946	9	3 0	0.9968	0.0196	18.0	0.75	40	26.7
947	9	30	0.9968	0.0196	21.1	0.75	40	28.0
948	ý	3 0	0.9968	0.0196	23,5	0.75	40	30.3
949	9	30	0.9968	0.0196	26.4	0,75	40	32,8
950	· ģ	30	0.9968	0.0196	29.0	0,75	40	40.0
951	9	30	0.9967	0.0197	11.9	1.00	40	17.8
952	9	30	0.9967	0.0197	15.1	1.00	40	22.0
953	9	30	0.9967	0.0197	18.0	1.00	40	23.0
954	9	3 0	0.9967	0.0197	21.1	1.00	40	27.0
955	9	3 0	0.9967	0.0197	23.5	1,00	40	27,2
956	9	30	0.9967	0.0197	26.4	1,00	40	32.3
957	9	3 0	0.9967	0.0197	29.0	1,00	40	37.0
960	- 9	30	0.9983	0.0100	16.6	0.10	40	40.0
961	9	30	0.9983	0.0100	16.6	0.31	40	33+0
962	9	30	0.9983	0.0100	25.2	0.31	40	23.0
963		30	0.9987	0.0100	16.6 25.2	0.50	40	40.0
964 965	9	30 30	0.9983	0.0100	33.0	0.50	40	40.0
966	9	. 3 0	0.9983	0.0100	40.8	0.50	40	40.0
967	- 	3 0	0,9983	0.0100	16.6	0,75	40	22.0
968	9	30	0.9983	0.0100	25.2	0.75	40	22.0
969	- 9	30	0.9983	0.0100	33.0	0.75	40	31.0
970	9	3 0	0.9983	0.0100	40.8	0.75	40	40.0
971	9	3 0	0.9982	0.0103	16.6	1,00	40	20.0
972	9	3 0	0.9982	0.0103	25.2	1,00	40	25+3
973	9	30	0.9982	0.0103	33.0	1,00	40	27.6
974	9	30	0.9982	0.0103	40.8	1.00	40	36.2
975	9	30	0.9982	0.0103	16.6	0.75	30	22.0
976 977	9	30	0.9982	0.0103	25.2 21.1	0,75	30 30	29.0 22.0
977 978	9	30 30	0.9982	0.0103	16.6	0,50	30	25.8
979	9	30	0.9982	0.0103	21.1	0.50	30	30.0
980	ģ	30	0.9982	0.0103	16,6	0.50	30	30.0
<u> </u>			0,7,7,0,0					···
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	<u>_</u>							
	<u> </u>		,					

(7 of 7 sheets)

Table 2 Dilution Tests

	D _o	D	Δρ _m		U	H				` Min	nimum O	bserve	d Dilu	tion (Dimens	ionles	s) at	Indica	ted Dis	tance	Downstr	eam fr	om Diff	ıser (f			_
Test	in.	in.	g/cc	$\mathbf{F}_{\mathtt{D}}$	knots	ft	5	10	<u>15</u>	50	_30	40	50	<u>60</u>	_80	100	120	160	200	220	240	280	320	400	440	480	500
420 421 615 617 619	3 6 6	20 20 20 20	0.021 0.022 0.021 0.021 0.021	36.3 36.3 13.1 18.9 25.0	0.5 1.0 0.1 0.1	40 40 40 40 40	ŶΩ	80 37 36	95	63 49	108 48	- 83		-	489 980		532	792	703 1619			-	1380 3388				
621 623 625 627 629	6 6 6 6	20 20 20 20	0.021 0.021 0.021 0.021 0.021	32.2 13.2 18.3 25.1 32.6	0.1 0.2 0.2 0.2 0.2	ちちちちち		37 64		42 21 20	37 83	78 44	75 108	81 41	103 55	68			-								
631 633 635 637 - 639	6 6 6 6	20 20 20 20 20	0.021 0.021 0.021 0.021 0.021	13.0 19.0 25.2 32.9 13.2	0.3 0.3 0.3 0.3 0.44	5 5						64 56 66 77 70	-i		124 85 87 102	142	190 95 137 158	215 153 131 160 156			197 200						
641 643 645 655 657	6 . 6 6 6	20 20 20 20 20	0.021 0.021 0.021 0.021 0.021	19.0 24.7 32.7 13.3 18.7	0.44 0.44 0.5 0.5	40 40 40 40					-	72 40 44 80 71			124 148	148 59 75	121 138	182 78 86 215 134	182 177	90	212	188					
659 661 663 665 667	6 6 6 6	20 20 20 20 20	0.021 0.021 0.021 0.021 0.021	24.6 32.1 13.1 18.6 24.2	0.5 0.5 0.75 0.75 0.75	4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4						85 64 157 125 128			172 115 197 240 143		132 94 219	148 137 285 302 209	170 426		168 177 436 408 387	14145	460 460	358	660 542		603
669 671 673 675 677	6 6 6 6	20 20 20 20 20	0.021 0.021 0.021 0.021 0.021	31.7 12.9 18.3 24.6 32.1	0.75 1.0 1.0 1.0	49 49 49 49						95 337 250 237 178			178 365 365 382 363			218 680 570 570 7 0 0			296 1060 840 715 940		396 1020 1170 900 980	396 2000 1420	1520	1780	572
686 694 936 950	6 6 9 9	20 20 30. 30	0.011 0.011 0.021 0.022	26.8 26.8 11.5 11.5	0.3 0.5 0.5 1.0	40 40 40						116 105 76			238	124	241 336	207 264			401		909 490	•			, -
4-615* 4-655* 4-659* 4-671* 4-675*	.6 6 6	20 20 20 20 20	0.021 0.021 0.021 0.021 0.021	13.1 13.1 24.8 13.1 24.8	0.1 0.5 0.5 1.0 1.0	5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5	33					150 81			207 184		•	1 11 128				348 231					
Heated** Outfall 1† Outfall 2† Outfall 3†	6 	20 20 20 20	0.021 0.021 0.021 0.021	18.8 	0.5 0.1 0.5 1.0	49 49 49 40						113			157		210					422			•		

^{* 4} ports at 13 ft o.c. ** Heated brine. † Q_b = 5.91 cfs.

Table 3
Lateral Spread Coordinates

					Co	ordin	ates	of La	teral	$_{f L}$ Spr ϵ	ead of	Bri	ne, f	`t				
<u>l'est</u>	x	<u>w</u>	х	W	x	w	х	w	х	w	<u>x</u>	W	x	w	<u>x</u>	w	<u>x</u>	V
07	40	33	80		100													
08 09	40	30 26	80 80	-	120 120	91	140	112										
10	40	23	80	48	120	86	470 470		150	113								
11	40	15	80	_	120		- 0											
12	20	22	80	36	120		140	104							'			
<u> 13</u>	40	25	80	23	120	40		101										
14	40	20	60	33	80	38	100		140	64	180	87						
15	40	11.	60	23	80		100	50	140		180	<u>87</u> 97						
16 17	40 40	15 17	60. 60	25 20	8 O 8 O	47	100 100	57 31	140		180 180	87						
18	40	16	60	22	80		100		140		180		220	97	260	107		
19	46	25	60	29	80		100		140		180	76	_	98				
20	40	23	60	30	80	37		38	140		180		220	96				
21	40	8	80	26	120		160		200	42	240		280	62				
22	40	14	80	24	120	2 5	160	44	500	49	240	51	280		320	65.		
23	40	16	80		120		160	35	200		240		280		320	75		
24 25	40 40	17 16	8 n 8 n	19	120 120	3 6 30	160 160	36 29	200	42	240 240	60 60	280 280		320 320	80 72		
<u> 25</u> 26	40	19	8 n		120		160	3 5	200	40	240	<u> 53</u>	280	65		30		
27	40	20	80		120		160	33	200	44	240	41	280		320	72		
28	40	9	80		120		160		200	23		30	280		320	44		
29	40	_12_	80	21	120	28	160	36	200	36	240		200		32n	54		
30	40	10	80		120	_	160	34	200		240		580		320	60		
31	40	19	80		120	26	160	29	200	25	240	46	280	48	320 320	72		
33 34	4 0 4 0	17 18	8 Ç 8 Q		120 120	33 25	160 160	. 37 34	200 200	43 51	240 240	48 48	280 280		320	64 52		
35	40	6	80	 7	120	13	160	57	120	24	240	23	<u> </u>	''	<u> </u>			
36	40	7	80		120		160	19	200	16	240		280	40				
37	40	9	80	1,8	120	28	160	25	200	25	240	35	240	37				
38	40	6	80		120		160	15	200	19	240		280	23				
3.9	40	12	80		120		160		200	30	240	34			320	40		
40 41	4 Ü	11 15	80		120	2 <u>4</u> 29	160 160	$\frac{29}{31}$	200	3 <u>1</u> 33	240	<u>33</u> 34	280 280	40	320 320	3 <u>2</u>		
42	40	8	80		120 120	18	160	20	200	18	240	17	280	20	220	77 44		
43	40	8	80	11	120	17	160		200	24	240	31						
44	40	8_	80		120		160	21	200	21								
45	40	10	80	14	120	2 5	160	27	200	25								
46	40	$\frac{11}{15}$	80		120	15	160	21	200		240	20	280	22				
47 48	4 0 4 0	15 15	80 80	23	120 120	35 25	160 160	35 33	200	36 32								
55	40	17	80		120	30	100											
56	40	18	80		120													
57	40	24	80	26	120	24	di barban andibid 1499	LATER STREET, SECTION OF	****						,		•	
58	40	12	80		120	17												
59 40	40	17	80		120	27												
60 62	40	19_	80	- <u>22</u> 18	120	27		·										
63	40	15	80	23														
64	40	17	80	$-\frac{23}{21}$	120	27												
23	40	59	60	110	_ = 0													
24	4 0	59	60	110														
25	40	20 19		114														
26	4 Ú		70	100														

Table 3 (Continued)

Test	x	w.	х.	W	x	Coord W	inate x	s of w	Late:	ral S w	pread x	of B	rine,	ft		7.7		
			_	_										<u>w</u>	<u>_x</u>	<u>w</u>	<u> </u>	W
27 28	4 0 4 0	20 36	80 80		100 120	118 118												
31	40	20	60	26	TCA	110			_									
32	40	19	60	33	80	62												
33	40	32	60	45	80		100	71	140	104							1	
34	40	23	6 p	28	80	35	100		140	105								
35	40	24	60	31	80		100	55	140		160							
36	40	28	60	33	80		100		140		180	95						
537 538	4 U 4 U	28 31	60	28 39	80 80		100	38	140		180	93	226	0.4			•	
39	40	17	60 80	27	120	62	160	86	140	20	180		220	70				
540	40	20	80		120		160		200	97			•					
541	40	16	80	30	120	52	160	87	200	102				-				
142	40_	15	80	27	120		160	78	200	94	,							
543	40	17	8 0	26	120	51	160	74	200	95								
544	40	20	80		120		160		200	91								
545	40	22	80		120		160		200	77	230	92						
54 <u>6</u> 547	40	- <u>14</u> 18	80 80		120 120		160 160	54	200	7 <u>3</u> 87		<u>97</u> 100						
548	40	17	80	26	120	45	160		200	91	270	100						
549	40	11	80	25	120	49	160	65	200		240	99						
550	40	17	80		120	45	160		200		240	92						
551	4 0	24	80	27			160	57	200	82								
552	40	26	80		120		160		200		240	90						
553	40	23	9,0	36	120		160	60	200	78	240	84	AU.	D 4				
554 555	40	_30 17	8 <u>0</u> 60	20	120 80		160		200 140	56	240 180	70	280	84				
556	40	18	.60	21	80		100	33	140		180	65 70	220	8 g 8 5				
557	40	21	60	23	80		100	36	140		180		220	81				
558	4.0	21	60	25	80		100	35			180	64	220	80				
559	40	2 6	60	29	80	30	100	35	140		180	70	550	83				
560	40	25	60	<u>30</u>	80		100	35	140		180	57	220		260	90		
661	40	21	60	24	80	28	100	36	140		180	60	220	74				_
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565	40	23	60	22	80		100	27	140		180	43	220	53	260		300	7
666	4 0	25	60	26	80		100	32			180	43	220		260		300	6
567	40	. 15	60	19	8 g	25	100	29	140		180	49	220		260	53	300	6
668	40	20	60	21	80		100	34	140		180	46	220		260		300	6.
669	40	20	60	23	80	2/	100	29	140		180	44	220		260		300	6
67 <u>0</u> 671	40	<u>19</u>	60 80	23 17	80 120	27	100 160	28 22	140 200	<u>35</u>	180 240	27	220	46	260 320		300	6
672	40	14	80	17			160		200		240	-	280	7 -	320	28 33		
673	40	21	80	28			160	39		41	240	41			320	42		
674	40	15	80	22	120	30	160	38	200	40	240		280	40	320	40		
675	40	10	80	15	120	22	160	31	200	41	240	44	280	47				
676	40	13	<u>8 p</u>		120		160		200		240	55	280	58	320	60		
677 670	40	23	80	37	-		160		200		240		280		320	50		
67 <u>8</u> 682	40	<u> 16</u> 29	80 80	67	120 120	111	160		200	40	240	58_	280	ა/	320	38	 ,	
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684	40	32	80		120		160		200	110	225	125						
	40	21.	80		120		160	85										

Table 3 (Continued)

					Coor	dinet	es of	Lete	ral	Spree	d of	Brine	, ft	· ;			
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	<u> </u>	<u> </u>							_				<u> </u>				
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687	4 0	20	80	27 120	32	160		200	86								
688	40	27	80	40 120				200	63	544							
689	40	20	80	26 120	38	160		200	65 75	240	81						
690 691	40	19 17	80 80	27 120 26 120	35	160 160		240	75	240	86						· ·
692	40	21	80	37 120		160	40		42		_	280	47	320	5 <i>7</i>		
693	40	18	80	26 120		160		200		240		260	76	920			
694	40	18	80	26 120		160		200	70	240	81	•	, -				
695	40	23	80	33 120	37	160		200	58	240	72	280	82				
696	40	24	80	29 120	28	160		200	35	240	40	280	44			_;	
697	40	13	80	20 120		160		200	32	240		280	45				
698	40	18	80	27 120		160		200	<u> </u>		47			320			
699	40	15	80	24 120		160		200		240		280		320	51		
700	40	19	80	30 120		160	36	200	40	240	_47	280	22	320	53		
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916	40	24	60	46 80	73										مه ، سر		
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924	40	18	80	26 120	36	160		200	96			-					
925	40	20	80	35 120				200		240	116						
926	40	29	80	39 120	40	160		200	70	240	101						
927	40	39	80	43 120	47	160		200	70	240		280	119				
928	40	46	8 C	64 120	80	160		200	83	240	114					× .	
930	40	21	80	30 120		160		200	93								
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936	60	29	80	31 100		140		180	91	_ , ,	•	- ∨					
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940		32	80	37 100		140		180	49	220	72	260	85	300	<u> 101</u>		
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942	60 40	36 19	<u>80</u> 60	42 100 24 80	25	140 100		180 140		220 180		260 220		300 260	<u>82</u>	300	67
944	40	23	60	26 80		100		140		180		220		260		300	77
945	40	21	6 p	26 80		100		140		160		220		260		300	83
946	40	19	60	24 80		100		140		180		220		260		300	82
947	40	24	60	27 80	30	100	33	140	50	180	58	220	69	260	78	300	88
948	40	25	60	32 80		100		140		180		220		260		300	87
949	40	26	60	32 80	37			140		180		220		260		300	49
950	40	17	80	22 120	24	160		200		240		280		320	54		
951	40	12	80	1 7 120	20	160	_	200		240	43	280	40	320	48		
							— (Co	ntin	æ d)		-			(:	3 of	4 s he	ets)

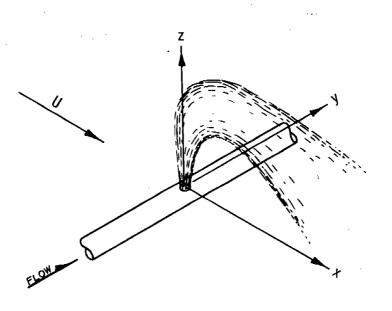
Table 3 (Concluded)

							Coo	rdina	tes c	f Lat	eral	Spre	ad of	Brin	e, ft	;			
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APPENDIX A: NOTATION

- A Cross-sectional area of outfall, ft²
- An Cross-sectional area of port, ft2
- B Dimensionless intercept
- B_{T_2} Calibration intercept at temperature T_2 , g/cc
- B_{T_3} Calibration intercept at temperature T_3 , g/cc
 - C Dimensionless coefficient
- C Conductivity meter voltage for a salt solution, mv
- D Outfall diameter, ft or in.
- D Port diameter, ft or in.
- $\mathbb{F}_{\mathbb{D}}$ Densimetric port Froude number, $\mathbb{V}\sqrt{\frac{\Delta \rho_{m}}{\rho_{f}}} gD_{o}$
 - g Gravitational acceleration, ft/sec2
- H Ambient flow depth, ft
- Q_h Brine discharge, cfs
 - R Exponent, see equation 9
- \mathbb{R}_c Channel Reynolds number
- R Port Reynolds number
- T Temperature, C
- T_a, T_b Temperatures of brine solutions, C
 - T₁ Temperature when solution densities are checked, C
 - T₂ Temperature when probes are calibrated, C
 - T_3 Temperature at which data are taken, C
 - U Ambient flow velocity, fps
 - V Average port velocity, fps
 - w Plume width, ft
 - w_0 Total plume width at x_0 , ft
 - x Distance downstream from port center line, ft
 - $\mathbf{x}_{_{\scriptsize{\textsc{O}}}}$ Downstream distance from port at which plume falls to bottom, ft
 - y Distance parallel to diffuser in reference to center of a given port, ft
 - z Elevation above bottom, ft

Maximum height of upper boundary of jet above bottom, ft Z_{m} Far-field effluent density minus ambient fluid density, g/cc Δρ Initial effluent density minus ambient fluid density, g/cc $\triangle P_{\mathbf{m}}$ $\triangle \rho_m / \triangle \rho = dilution$ ϵ_{m} Minimum observed dilution Kinematic viscosity of water, ft²/sec Density, g/cc Densities of brine solutions, g/cc ρ_a, ρ_b Ambient fluid density, g/cc ρ_{f} Solution density at temperature T_1 , g/cc ρ_{sl} Solution density at temperature T_2 , g/ccDensities of distilled water at temperatures T_1, T_2, T_3 , g/ccρ₁,ρ₂,ρ₃ Indicates functional relation



Definition of coordinate system

APPENDIX B: CONDUCTIVITY PROBE CALIBRATION AND DATA REDUCTION

Density is assumed to vary as a function of temperature and salinity only, and conductivity is assumed to vary only with salinity over the small range of temperatures encountered in the testing. At a given temperature, density varies linearly with conductivity (salinity); and at a different temperature, the linear variation has the same slope but is displaced by the difference in the density of distilled water at the two temperatures. An equivalent statement is that the nonlinear temperature-density relation for distilled water is linearly displaced upward by small changes in density due to salinity. Fig. Bl illustrates these relations.

In practice, the temperature of the calibrating solutions differed from one to another; the temperature at which the solutions were checked for density was different from that at which the probes were calibrated; and temperatures during actual testing were still different. In a FORTRAN program written to reduce the calibrations and compute dilutions, the following equation for the density of distilled water is used to make temperature adjustments:

$$\rho = 1 - \left[\frac{(T - 3.9863)^2}{508,929.2} \cdot \frac{T + 288.9414}{T + 68.12963} \right]$$
 (B1)

where

 ρ = density, g/cc

T = temperature, C

Letting

 ρ_{sl} = solution density at temperature T_l

 ρ_{s2} = solution density at temperature T_2

 ρ_1 = density of distilled water at temperature T_1

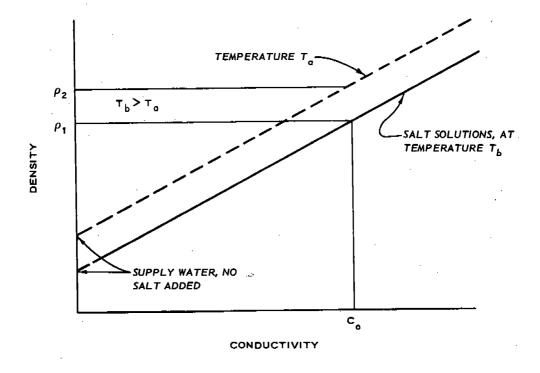
 ρ_2 = density of distilled water at temperature T_2

 \boldsymbol{T}_{1} = temperature when solution densities are checked

 T_2 = temperature when probes are calibrated

the solution densities are converted to density at the calibration temperature by the following equation:

^{*} See Literature Cited at end of main text, p 38.



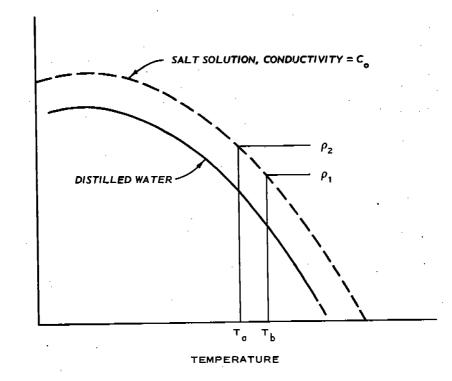


Fig. Bl. Conductivity-temperature-density relations

$$\rho_{s2} = \rho_{s1} + (\rho_1 - \rho_2)$$
 (B2)

At the calibration temperature T_2 , linear calibrations of density versus conductivity can be constructed for each probe. A least-squares slope is computed and the intercept calculated by forcing the calibration to pass through the first calibration point, which in practice was "fresh" water; this procedure was necessary to ensure that calibration error was minimized at the sensitive lower ranges of conductivity. These intercepts, representing the density of water having zero conductivity (but denser than distilled water due to suspended solids, etc.), are functions only of temperature. Thus the calibrations of temperature T_2 are adjusted to temperature T_3 by

$$B_{T_3} = B_{T_2} + (\rho_2 - \rho_3)$$
 (B3)

where

 B_{T_2} = calibration intercept at temperature T_2 B_{T_3} = calibration intercept at temperature T_3 p_3 = density of distilled water at temperature T_3 T_3 = temperature at which data are taken

In reducing the raw data, the FORTRAN program first processes the calibration data as described above to yield basic linear calibrations. The raw data input consists of grouped conductivities and temperatures, along with a "background" conductivity that represents the approximate freshwater conductivity during the test for each probe. Although the density of the background remains essentially constant at a given temperature, its conductivity may shift slightly from its calibration value due to small changes in background salinity. Therefore, each conductivity reading is slightly adjusted by an amount equal to the background shift between the time of calibration and time of test so that calibrations are continuously updated.

For a given data point, the calibration for that probe is shifted to the data temperature T_3 , and a density is computed. In computing dilution, the initial density difference $\Delta_{\!\rho_m}$ is calculated from hydrometer readings for the brine and ambient fluid, adjusted to the same temperature. The diluted density difference $\Delta_{\!\rho}$ is the difference between the density as computed from the temperature-conductivity data and the density of the ambient fluid, adjusted to the data temperature. Dilution is then calculated by

 $\epsilon = \frac{\Delta \rho_{\rm m}}{\Delta_0} \tag{B4}$